# **Classical TGD**

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#### Abstract

This chapter is devoted to the applications of what might be called classical TGD to astrophysics and cosmology. In a well-defined sense classical TGD defined as the dynamics of space-time surface determining them as kind of generalized Bohr orbits can be regarded as an exact part of quantum theory and assuming quantum classical correspondence has served as an extremely valuable guideline in the attempts to interpret TGD, to form a view about what TGD really predicts, and to to guess what the underlying quantum theory could be and how it deviates from standard quantum theory. Also TGD inspired cosmology and astrophysics relies on this general picture.

#### 1. Many-sheeted cosmology

The many-sheeted space-time concept, the new view about the relationship between inertial and gravitational four-momenta, the basic properties of the paired cosmic strings, the existence of the limiting temperature, the assumption about the existence of the vapor phase dominated by cosmic strings, and quantum criticality imply a rather detailed picture of the cosmic evolution, which differs from that provided by the standard cosmology in several respects but has also strong resemblances with inflationary scenario.

The most important differences are following.

a) Many-sheetedness implies cosmologies inside cosmologies Russian doll like structure with a spectrum of Hubble constants.

b) TGD cosmology is also genuinely quantal: each quantum jump in principle recreates each sub-cosmology in 4-dimensional sense: this makes possible a genuine evolution in cosmological length scales so that the use of anthropic principle to explain why fundamental constants are tuned for life is not necessary.

c) The new view about energy means that inertial energy is negative for space-time sheets with negative time orientation and that the density of inertial energy vanishes in cosmological length scales. Therefore any cosmology is in principle creatable from vacuum and the problem of initial values of cosmology disappears. The density of matter near the initial moment is dominated by cosmic strings approaches to zero so that big bang is transformed to a silent whisper amplified to a relatively big bang.

d) Dark matter hierarchy with dynamical quantized Planck constant implies the presence of dark space-time sheets which differ from non-dark ones in that they define multiple coverings of  $M^4$ . Quantum coherence of dark matter in the length scale of space-time sheet involved implies that even in cosmological length scales Universe is more like a living organism than a thermal soup of particles.

e) Sub-critical and over-critical Robertson-Walker cosmologies are fixed completely from the imbeddability requirement apart from a single parameter characterizing the duration of the period after which transition to sub-critical cosmology necessarily occurs. The fluctuations of the microwave background reflect the quantum criticality of the critical period rather than amplification of primordial fluctuations by exponential expansion. This and also the finite size of the space-time sheets predicts deviations from the standard cosmology.

#### 2. Cosmic strings

Cosmic strings belong to the basic extremals of the Kähler action. The string tension of the cosmic strings is  $T \simeq .2 \times 10^{-6}/G$  and slightly smaller than the string tension of the GUT strings and this makes them very interesting cosmologically. Concerning the understanding of cosmic strings a decisive breakthrough came through the identification of gravitational four-momentum as the difference of inertial momenta associated with matter and antimatter and the realization that the net inertial energy of the Universe vanishes. This forced to conclude cosmological constant in TGD Universe is non-vanishing. p-Adic length fractality predicts that  $\Lambda$  scales as  $1/L^2(k)$  as a function of the p-adic scale characterizing the space-time sheet. The recent value of the cosmological constant comes out correctly. The gravitational energy density described by the cosmological constant is identifiable as that associated with topologically condensed cosmic strings and of magnetic flux tubes to which they are gradually transformed during cosmological evolution. p-Adic fractality and simple quantitative observations lead to the hypothesis that pairs of cosmic strings are responsible for the evolution of astrophysical structures in a very wide length scale range. Large voids with size of order 10<sup>8</sup> light years can be seen as structures containing knotted and linked cosmic string pairs wound around the boundaries of the void. Galaxies correspond to same structure with smaller size and linked around the supra-galactic strings. This conforms with the finding that galaxies tend to be grouped along linear structures. Simple quantitative estimates show that even stars and planets could be seen as fractal cosmic necklace consisting of cosmic strings linked like pearls around longer cosmic strings linked like...

#### 3. Dark matter and quantization of gravitational Planck constant

The notion of gravitational Planck constant having gigantic value is perhaps the most radical idea related to the astrophysical applications of TGD. D. Da Rocha and Laurent Nottale have proposed that Schrödinger equation with Planck constant  $\hbar$  replaced with what might be called gravitational Planck constant  $\hbar_{gr} = \frac{GmM}{v_0}$  ( $\hbar = c = 1$ ).  $v_0$  is a velocity parameter having the value  $v_0 = 144.7 \pm .7$  km/s giving  $v_0/c = 4.6 \times 10^{-4}$ . This is rather near to the peak orbital velocity of stars in galactic halos. Also subharmonics and harmonics of  $v_0$  seem to appear. The support for the hypothesis coming from empirical data is impressive.

Nottale and Da Rocha believe that their Schrödinger equation results from a fractal hydrodynamics. Many-sheeted space-time however suggests astrophysical systems are not only quantum systems at larger space-time sheets but correspond to a gigantic value of gravitational Planck constant. The gravitational (ordinary) Schrödinger equation would provide a solution of the black hole collapse (IR catastrophe) problem encountered at the classical level. The resolution of the problem inspired by TGD inspired theory of living matter is that it is the dark matter at larger space-time sheets which is quantum coherent in the required time scale.

I have proposed already earlier the possibility that Planck constant is quantized and the spectrum is given in terms of logarithms of Beraha numbers: the lowest Beraha number  $B_3$  is completely exceptional in that it predicts infinite value of Planck constant. The inverse of the gravitational Planck constant could correspond a gravitational perturbation of this as  $1/\hbar_{gr} = v_0/GMm$ . The general philosophy would be that when the quantum system would become non-perturbative, a phase transition increasing the value of  $\hbar$  occurs to preserve the perturbative character and at the transition  $n = 4 \rightarrow 3$  only the small perturbative correction to  $1/\hbar(3) = 0$  remains. This would apply to QCD and to atoms with Z > 137 as well.

TGD predicts correctly the value of the parameter  $v_0$  assuming that cosmic strings and their decay remnants are responsible for the dark matter. The harmonics of  $v_0$  can be understood as corresponding to perturbations replacing cosmic strings with their n-branched coverings so that tension becomes  $n^2$ -fold: much like the replacement of a closed orbit with an orbit closing only after n turns. 1/n-sub-harmonic would result when a magnetic flux tube split into n disjoint magnetic flux tubes. An attractive solution of the matter antimatter asymmetry is based on the identification of also antimatter as dark matter.

# 1 Introduction

A brief summary of what might be called basic principles is in order to facilitate the reader to assimilate the basic tools and rules of intuitive thinking involved.

#### 1.1 Quantum-classical correspondence

The fundamental meta level guiding principle is quantum-classical correspondence (classical physics is an exact part of quantum TGD). The principle states that all quantum aspects of the theory, which means also various aspects of consciousness such as volition, cognition, and intentionality, should have space-time correlates [10]. Real space-time sheets provide kind of symbolic representations whereas p-adic space-time sheets provide correlates for cognition and intentions. All that we can symbolically communicate about conscious experience relies on quantal space-time engineering to build these representations.

#### **1.2** Classical physics as exact part of quantum theory

Classical physics corresponds to the dynamics of space-time surfaces determined by the absolute minimization of Kähler action or some other principle selecting preferred extremals of Kähler action [E2]. This dynamics have several unconventional features basically due to the possibility to interpret the Kähler action as a Maxwell action expressible in terms of the induced metric defining classical gravitational field and induced Kähler form defining a non-linear Maxwell field not as such identifiable as electromagnetic field however.

# 1.2.1 Classical long ranged weak and color fields as signature for a fractal hierarchy of copies standard model physics

The geometrization of classical fields means that various classical fields are expressible in terms of imbedding space-coordinates and are thus not primary dynamical variables. This predicts the presence of long range weak and color (gluon) fields not possible in standard physics context. It took 26 years to end up with a convincing interpretation for this puzzling prediction.

What seems to be the correct interpretation is in terms of an infinite fractal hierarchy of copies of standard models physics with appropriately scaled down mass spectra for quarks, leptons, and gauge bosons. Both p-adic length scales and the values of Planck constant predicted by TGD [C6] label various physics in this hierarchy. Also other quantum numbers are predicted as labels. This means that universe would be analogous to an inverted Mandelbrot fractal with each bird's eye of view revealing new long length scale structures serving also as correlates for higher levels of self hierarchy.

Exotic dark weak forces and their dark variants are consistent with the experimental widths for ordinary weak gauge bosons since the particles belonging to different levels of the hierarchy do not have direct couplings at Feynman diagram level although they have indirect classical interactions and also the de-coherence reducing the value of  $\hbar$  is possible. Classical long ranged weak fields play a key role in quantum control and communications in living matter [M3, N4]. Long ranged classical color force in turn is the backbone in the model of color vision [K3]: colors correspond to the increments of color quantum numbers in this model. The increments of weak isospin in turn could define the basic color like quale associated with hearing (black-white  $\leftrightarrow$  to silence-sound [K3, M5, M6]).

#### **1.2.2** Topological field quantization and the notion of many-sheeted space-time

The compactness of  $CP_2$  implies the notions of many-sheeted space-time and topological field quantization. Topological field quantization means that various classical field configurations decompose into topological field quanta. One can see space-time as a gigantic Feynman diagram with lines thickened to 4-surfaces. Absolute minimization of Kähler action (or any other analogous selection principle) implies that only selected field configurations analogous to Bohr's orbits are realized physically so that quantum-classical correspondence becomes very predictive. An interpretation as a 4-D quantum hologram is a further very useful picture [K2] but will not be discussed in this chapter in any detail.

Topological field quantization implies that the field patterns associated with material objects form extremely complex topological structures which can be said to belong to the material objects. The notion of field body, in particular magnetic body, typically much larger than the material system, differentiates between TGD and Maxwell's electrodynamics, and has turned out to be of fundamental importance in the TGD inspired theory of consciousness. One can say that field body provides an abstract representation of the material body.

One implication of many-sheetedness is the possibility of macroscopic quantum coherence. By quantum classical correspondence large space-time sheets as quantum coherence regions are macroscopic quantum systems and therefore ideal sites of the quantum control in living matter.

- 1. The original argument was that each space-time sheet carrying matter has a temperature determined by its size and the mass of the particles residing at it via de Broglie wave length  $\lambda_{dB} = \sqrt{2mE}$  assumed to define the p-adic length scale by the condition  $L(k) < \lambda_{dB} < L(k_{>})$ . This would give very low temperatures when the size of the space-time sheet becomes large enough. The original belief indeed was that the large space-time sheets can be very cold because they are not in thermal equilibrium with the smaller space-time sheets at higher temperature.
- 2. The assumption about thermal isolation is not needed if one accepts the possibility that Planck constant is dynamical and quantized and that dark matter corresponds to a hierarchy of phases characterized by increasing values of Planck constant [C6, J6]. From E = hfrelationship it is clear that arbitrarily low frequency dark photons (say EEG photons) can have energies above thermal energy which would explain the correlation of EEG with consciousness. This vision allows to formulate more precisely the basic notions of TGD inspired theory of consciousness and leads to a model of living matter giving precise quantitative predictions. Also the ability of this vision to generate new insights to quantum biology provides strong support for it [M3].

Many-sheeted space-time predicts also fundamental mechanisms of metabolism based on the dropping of particles between space-time sheets with an ensuing liberation of the quantized zero point kinetic energy. Also the notion of many-sheeted laser follows naturally and population inverted many-sheeted lasers serve as storages of metabolic energy [K6].

Space-time sheets topologically condense to larger space-time sheets by wormhole contacts which have Euclidian signature of metric. This implies causal horizon (or elementary particle horizon) at which the signature of the induced metric changes from Minkowskian to Euclidian. This forces to modify the notion of sub-system. What is new is that two systems represented by space-time sheets can be unentangled although their sub-systems bound state entangle with the mediation of the join along boundaries bonds connecting the boundaries of sub-system space-time sheets. This is not allowed by the notion of sub-system in ordinary quantum mechanics. This notion in turn implies the central concept of fusion and sharing of mental images by entanglement [10].

#### 1.2.3 The possibility of negative energies

A further prediction derives from the fact that space-time is 4-surface rather than an abstract manifold. The energy momentum tensor of general relativity is replaced by a collection of conserved energy and momentum currents, which are 4-vector fields. This makes the notions of energy and momentum precisely defined but also implies that the sign of energy and momentum depend on the time-orientation of the space-time sheet. Negative energies become therefore possible somewhat like in the lines of a Feynman diagram. Negative energy topological light rays have phase conjugate laser waves [24] as the most plausible standard physics counterparts, and play a fundamental role in quantum metabolism as a kind of quantum credit card [K6]. They generate also time like entanglement which corresponds to a formation of new kind of bound states.

Negative energies might be possible even for ordinary particles and could mean dramatic deviation from the standard quantum theory. The roles of annihilation and creation operators have changed for negative energy space-time sheets. This would mean that operator combinations involving both annihilation and creation operators would generate states involving positive and negative energy space-time sheets. One can even imagine that a intentional action could create states with vanishing net quantum numbers and that positive and negative energy particles could be separated from each other.

#### 1.2.4 TGD Universe is quantum spin glass

Since Kähler action is Maxwell action with Maxwell field and induced metric expressed in terms of  $M_+^4 \times CP_2$  coordinates, the gauge invariance of Maxwell action as a symmetry of the vacuum extremals (this implies is a gigantic vacuum degeneracy) but not of non-vacuum extremals. Gauge symmetry related space-time surfaces are not physically equivalent and gauge degeneracy transforms to a huge spin glass degeneracy. Spin glass degeneracy provides a universal mechanism of macro-temporal quantum coherence and predicts degrees of freedom called zero modes not possible in quantum field theories describing particles as point-like objects. Zero modes not contributing to the configuration space line element are identifiable as effectively classical variables characterizing the size and shape of the 3-surface as well as the induced Kähler field. Spin glass degeneracy as mechanism of macroscopic quantum coherence should be equivalent with dark matter hierarchy as a source of the coherence [K2].

#### 1.2.5 Classical and p-adic non-determinism

The vacuum degeneracy of Kähler action implies classical non-determinism, which means that space-like 3-surface is not enough to fix the space-time surface associated with it uniquely as an absolute minimum of action, and one must generalize the notion of 3-surface by allowing sequences of 3-surfaces with time like separations to achieve determinism in a generalized sense. These "association sequences" can be seen as symbolic representations for the sequences of quantum jumps defining selves and thus for contents of consciousness. Not only speech and written language define symbolic representations but all real space-time sheets of the space-time surfaces can be seen in a very general sense as symbolic representations of not only quantum states but also of quantum jump sequences. An important implication of the classical non-determinism is the possibility to have conscious experiences with contents localized with respect to geometric time. Without this non-determinism conscious experience would have no correlates localized at space-time surface, and there would be no psychological time.

p-Adic non-determinism follows from the inherent non-determinism of p-adic differential equations for any action principle and is due to the fact that integration constants, which by definition are functions with vanishing derivatives, are not constants but functions of the pinary cutoffs  $x_N$ defined as  $x = \sum_k x_k p^k \rightarrow x_N = \sum_{k < N} x_k p^k$  of the arguments of the function. In p-adic topology one can therefore fix the behavior of the space-time surface at discrete set of space-time points *above* some length scale defined by p-adic concept of nearness by fixing the integration constants. In the real context this corresponds to the fixing the behavior *below* some time/length scales since points p-adically near to each other are in real sense faraway. This is a natural correlate for the possibility to plan the behavior and p-adic non-determinism is assumed to be a classical correlate for the non-determinism of intentionality, and perhaps also imagination and cognition.

These two non-determinisms allow to understand the self-referentiality of consciousness at a very general level. In a given quantum jump a space-time surface can be created with the property that it represents symbolically or cognitively something about the contents of consciousness before the quantum jump. Thus it becomes possible to become conscious about being conscious of something. This is very much like mathematician expressing her thoughts as symbol sequences which provides feedback to go the next abstraction level.

Classical and p-adic non-determinisms force also the generalization of the notion of quantum entanglement. Time-like entanglement, crucial for understanding long term memory and precognition becomes possible. The notion of many-sheeted space-time forces also to modify the notion of sub-system, which implies that unentangled systems can have entangled sub-systems. One can partially understand this in terms of length scale dependent notion of entanglement (the entanglement of sub-systems is not seen in the length scale resolution defined by the size of unentangled systems) but only partially. The formation of join along boundaries bonds between sub-system space-time sheets and the fact that topologically condensed space-time sheets are separated by elementary particle horizons from larger space-time sheets, provide the deeper topological motivation for the generalization of sub-system concept.

#### 1.2.6 p-Adic fractality of life and consciousness

p-Adic fractality of biology and consciousness has become an increasingly important guide line in the construction of the theory. This notion allows to relate phenomena occurring in the molecular level to phenomena like remote viewing and psychokinesis and it leads also to the view that topological field quanta of various fields of astrophysical size are crucial for the functioning of bio-systems. If one accepts p-adic fractality, the theory can be tested in unexpected manners, in particular in molecular and cellular length scales where the systems are much simpler. Sensory perception, long term memory, remote mental interactions, metabolism: all these phenomena rely on the same basic mechanisms. p-Adic length scale hypothesis allows to quantify the hypothesis with testable quantitative predictions.

#### 1.2.7 Double slit experiment and classical non-determinism

Bohr's complementarity principle is the basic element of Copenhagen interpretation and at the same time one of the most poorly defined aspects of this interpretation. If the possibility of macroscopic quantum entanglement between measurement instrument and quantum system is accepted, complementary principle becomes un-necessary. This is however not all that is needed. If classical non-determinism makes it possible to represent quantum jump sequences at space-time level, a revision of space-time description of quantum measurement is necessary. This sounds very logical but to be honest, I write these lines only after having learned about the remarkable experiment done by Shahriar Afshar [28].

The variant of double slit experiment by Shahriar Afshar seems to contradict the Copenhagen interpretation which states that the particle and field aspects are complementarity and thus mutually exclusive. In the case of double slit experiment complementarity predicts that the measurement of whether the photon came to the detector through slit 1 or 2 should destroy the interference pattern of electromagnetic fields in the region behind the screen.

The experimental arrangement of Afshar differs from the standard double slit experiment in that a lens was added behind the screen. The lens transmitted the photons coming from slits 1 and 2 via mirrors to detectors A and B so that in particle picture a photon detected by A (B) could be regarded as coming from slit 1 (2). In the first step both slits were open and the detectors represented interference patterns representing diffraction through single slit. The other slit was then closed and metal wires at the positions of dark interference rings were added. These wires degraded somewhat the image in the second detector. After this the second slit was opened again. Surprisingly, the resulting interference pattern was the original one.

The measurement certainly measures the particle aspect of photons. On the other hand, the preservation of the detected patterns means that no photons did enter in the regions containing the wires so that also interference pattern is there. Hence wave and particle aspects seem to be mutually consistent. This finding is difficult to understand in Copenhagen interpretation and also in the manyworlds interpretation of quantum mechanics. Afshar himself suggest that the very notion of photon must be questioned. It is however difficult to accept this view since the photon absorption quite concretely corresponds to a click in the detector and also because the mathematical formalism of second quantization works so fantastically.

The conclusion can be criticized. What is primarily measured is not basically through which slit the photons came but whether the direction of the momentum of the photon emerging from the lens was in the angle range characterizing the detector or not. One can however argue that in deterministic physics for fields the two measurements are equivalent so that the problem remains.

In TGD framework the classical physics is not completely deterministic and this has led to a generalization of the notion of quantum classical correspondence. Space-time surface provides a classical (unfaithful) representation not only for quantum states but for quantum jump sequences or equivalently, for sequences of quantum states. The most obvious identification for the quantum states is as the maximal non-deterministic regions of a given space-time sheet.

In the recent context this would mean that the fields in the region between the screen and lens represent the state before the state function reduction and thus the interference pattern, whereas the fields in the region between lens and detectors represent the situation after the state function reduction. The interaction with lens involves classical non-determinism.

This picture conforms also with the notion of topological field quantization. The space-time decomposes into space-time sheets interpreted, topological field quanta (topological light rays containing photons, flux quanta of magnetic field, etc..). Topological field quanta correspond to the coherence regions for classical fields with spinor fields included. De-coherence corresponds to the splitting of space-time sheet to smaller, possibly parallel space-time sheets. Topological field quantum carries classical fields inside it but behaves as a whole like particle. Hence particle and wave aspects are consistent in the sense that below the size scale L of the topological field quantum (say the thickness of a magnetic flux tube or topological light ray) the description as a wave applies and above L particle description makes sense. In the recent case the coherence is lost at the lens space-time sheet where the space-time sheet representing interference pattern decomposes to two sheets representing photon beams going to the two detectors.

## 1.3 Some basic ideas of TGD inspired theory of consciousness and quantum biology

The following ideas of TGD inspired theory of consciousness and of quantum biology are the most relevant ones for what will follow.

1. "Everything is conscious and consciousness can be only lost" is the briefest manner to summarize TGD inspired theory of consciousness. Quantum jump as moment of consciousness and the notion of self are key concepts of the theory. Self is a system able to avoid bound state entanglement with environment and can be formally seen as an ensemble of quantum jumps. The contents of consciousness of self are defined by the averaged increments of quantum numbers and zero modes (sensory and geometric qualia). Moment of consciousness can be said to be the counterpart of elementary particle and self the counterpart of many-particle state, either bound and free. The selves formed by macro-temporal quantum coherence are in turn the counterparts of atoms, molecules and larger structures. Macro-temporal quantum coherence effectively binds a sequence of quantum jumps to a single quantum jump as far as conscious experience is considered. The idea that conscious experience is about changes amplified to macroscopic quantum phase transitions, is the key philosophical guideline in the construction of various models, such as the model of qualia, the capacitor model of sensory receptor, the model of cognitive representations, and declarative memories.

- 2. Macro-temporal quantum coherence is a second consequence of the spin glass degeneracy [K2]. It is essentially due to the formation of bound states and has as a topological correlate the formation of join along boundaries bonds connecting the boundaries of the component systems. During macro-temporal coherence quantum jumps integrate effectively to single long-lasting quantum jump and one can say that system is in a state of oneness, eternal now, outside time. Macro-temporal quantum coherence makes possible stable non-entropic mental images. Negative energy MEs are one particular mechanism making possible macro-temporal quantum coherence via the formation of bound states, and remote metabolism and sharing of mental images are other facets of this mechanism. The real understanding of the origin of macroscopic quantum coherence requires the generalization of quantum theory allowing dynamical and quantized Planck constant [J6, M3].
- 3. p-Adic physics as physics of intentionality and of cognition is a further key idea of TGD inspired theory of consciousness. p-Adic space-time sheets as correlates for intentions and p-adic-to-real transformations of them as correlates for the transformation of intentions to actions allow deeper understanding of also psychological time as a front of p-adic-to-real transition propagating to the direction of the geometric future. Negative energy MEs are absolutely essential for the understanding of how precisely targeted intentionality is realized.

# 2 Many-sheeted space-time, magnetic flux quanta, electrets and MEs

TGD inspired theory of consciousness and of living matter relies on space-time sheets carrying ordinary matter, topological light rays (massless extremals, MEs), and magnetic and electric flux quanta. There are some new results which motivate a separate discussion of them.

### 2.1 Dynamical quantized Planck constant and dark matter hierarchy

By quantum classical correspondence space-time sheets can be identified as quantum coherence regions. Hence the fact that they have all possible size scales more or less unavoidably implies that Planck constant must be quantized and have arbitrarily large values. If one accepts this then also the idea about dark matter as a macroscopic quantum phase characterized by an arbitrarily large value of Planck constant emerges naturally as does also the interpretation for the long ranged classical electro-weak and color fields predicted by TGD. Rather seldom the evolution of ideas follows simple linear logic, and this was the case also now. In any case, this vision represents the fifth, relatively new thread in the evolution of TGD and the ideas involved are still evolving.

#### 2.1.1 Dark matter as large $\hbar$ phase

D. Da Rocha and Laurent Nottale have proposed that Schrödinger equation with Planck constant  $\hbar$  replaced with what might be called gravitational Planck constant  $\hbar_{gr} = \frac{GmM}{v_0}$  ( $\hbar = c = 1$ ).  $v_0$  is a velocity parameter having the value  $v_0 = 144.7 \pm .7$  km/s giving  $v_0/c = 4.6 \times 10^{-4}$ . This is rather near to the peak orbital velocity of stars in galactic halos. Also sub-harmonics and harmonics of  $v_0$  seem to appear. The support for the hypothesis coming from empirical data is impressive.

Nottale and Da Rocha believe that their Schrödinger equation results from a fractal hydrodynamics. Many-sheeted space-time however suggests astrophysical systems are not only quantum systems at larger space-time sheets but correspond to a gigantic value of gravitational Planck constant. The gravitational (ordinary) Schrödinger equation would provide a solution of the black hole collapse (IR catastrophe) problem encountered at the classical level. The resolution of the problem inspired by TGD inspired theory of living matter is that it is the dark matter at larger space-time sheets which is quantum coherent in the required time scale [D6].

#### 2.1.2 Dark matter as a source of long ranged weak and color fields

Long ranged classical electro-weak and color gauge fields are unavoidable in TGD framework. The smallness of the parity breaking effects in hadronic, nuclear, and atomic length scales does not however seem to allow long ranged electro-weak gauge fields. The problem disappears if long ranged classical electro-weak gauge fields are identified as space-time correlates for massless gauge fields created by dark matter. Also scaled up variants of ordinary electro-weak particle spectra are possible. The identification explains chiral selection in living matter and unbroken  $U(2)_{ew}$  invariance and free color in bio length scales become characteristics of living matter and of biochemistry and bio-nuclear physics. An attractive solution of the matter antimatter asymmetry is based on the identification of also antimatter as dark matter.

#### 2.1.3 Dark matter hierarchy and consciousness

The emergence of the vision about dark matter hierarchy has meant a revolution in TGD inspired theory of consciousness. Dark matter hierarchy means also a hierarchy of long term memories with the span of the memory identifiable as a typical geometric duration of moment of consciousness at the highest level of dark matter hierarchy associated with given self so that even human life cycle represents at this highest level single moment of consciousness.

Dark matter hierarchy leads to detailed quantitative view about quantum biology with several testable predictions [M3]. The applications to living matter suggests that the basic hierarchy corresponds to a hierarchy of Planck constants coming as  $\hbar(k) = \lambda^k(p)\hbar_0$ ,  $\lambda \simeq 2^{11}$  for  $p = 2^{127-1}$ ,  $k = 0, 1, 2, \dots$  [M3]. Also integer valued sub-harmonics and integer valued sub-harmonics of  $\lambda$  might be possible. Each p-adic length scale corresponds to this kind of hierarchy and number theoretical arguments suggest a general formula for the allowed values of Planck constant  $\lambda$  depending logarithmically on p-adic prime [C6]. Also the value of  $\hbar_0$  has spectrum characterized by Beraha numbers  $B_n = 4\cos^2(\pi/n), n \geq 3$ , varying by a factor in the range n > 3 [C6].

The general prediction is that Universe is a kind of inverted Mandelbrot fractal for which each bird's eye of view reveals new structures in long length and time scales representing scaled down copies of standard physics and their dark variants. These structures would correspond to higher levels in self hierarchy. This prediction is consistent with the belief that 75 per cent of matter in the universe is dark.

#### 1. Living matter and dark matter

Living matter as ordinary matter quantum controlled by the dark matter hierarchy has turned out to be a particularly successful idea. The hypothesis has led to models for EEG predicting correctly the band structure and even individual resonance bands and also generalizing the notion of EEG [M3]. Also a generalization of the notion of genetic code emerges resolving the paradoxes related to the standard dogma [L2, M3]. A particularly fascinating implication is the possibility to identify great leaps in evolution as phase transitions in which new higher level of dark matter emerges [M3].

It seems safe to conclude that the dark matter hierarchy with levels labelled by the values of Planck constants explains the macroscopic and macro-temporal quantum coherence naturally. That this explanation is consistent with the explanation based on spin glass degeneracy is suggested by following observations. First, the argument supporting spin glass degeneracy as an explanation of the macro-temporal quantum coherence does not involve the value of  $\hbar$  at all. Secondly, the failure of the perturbation theory assumed to lead to the increase of Planck constant and formation of macroscopic quantum phases could be precisely due to the emergence of a large number of new

degrees of freedom due to spin glass degeneracy. Thirdly, the phase transition increasing Planck constant has concrete topological interpretation in terms of many-sheeted space-time consistent with the spin glass degeneracy.

#### 2. Dark matter hierarchy and the notion of self

The vision about dark matter hierarchy leads to a more refined view about self hierarchy and hierarchy of moments of consciousness [J6, M3]. The larger the value of Planck constant, the longer the subjectively experienced duration and the average geometric duration  $T(k) \propto \lambda^k$  of the quantum jump.

Dark matter hierarchy suggests also a slight modification of the notion of self. Each self involves a hierarchy of dark matter levels, and one is led to ask whether the highest level in this hierarchy corresponds to a single quantum jump rather than a sequence of quantum jumps. The averaging of conscious experience over quantum jumps would occur only for sub-selves at lower levels of dark matter hierarchy and these mental images would be ordered, and single moment of consciousness would be experienced as a history of events. One can ask whether even entire life cycle could be regarded as a single quantum jump at the highest level so that consciousness would not be completely lost even during deep sleep. This would allow to understand why we seem to know directly that this biological body of mine existed yesterday.

The fact that we can remember phone numbers with 5 to 9 digits supports the view that self corresponds at the highest dark matter level to single moment of consciousness. Self would experience the average over the sequence of moments of consciousness associated with each sub-self but there would be no averaging over the separate mental images of this kind, be their parallel or serial. These mental images correspond to sub-selves having shorter wake-up periods than self and would be experienced as being time ordered. Hence the digits in the phone number are experienced as separate mental images and ordered with respect to experienced time.

### 2.2 p-Adic length scale hypothesis and the connection between thermal de Broglie wave length and size of the space-time sheet

Also real space-time sheets are assumed to be characterized by p-adic prime p and assumed to have a size determined by primary p-adic length scale  $L_p$  or possibly n-ary p-adic length scale  $L_p(n)$ . Since multi-p-fractality is allowed [?], one cannot exclude even the possibility that each space-time dimension might correspond to its own p-adic length scale and even several p-adic primes could be associated with single dimension.

The possibility to assign a p-adic prime to the real space-time sheets is required by the success of the elementary particle mass calculations and various applications of the p-adic length scale hypothesis. Rational numbers are common to reals and all p-adic number fields. The p-adic-toreal transition transforming intentions to actions is made possible by a large number of common rational points between p-adic and real space-time surfaces, which supports the view that real space-time sheets obeys effective p-adic topology as an approximate topology in some resolution and below some length scale. p-Adic prime thus characterizes the classical non-determinism of the Kähler action.

Parallel space-time sheets with distance about  $10^4$  Planck lengths form a hierarchy. Each material object (...,atom, molecule, ..., cell,...) would correspond to this kind of space-time sheet. The p-adic primes  $p \simeq 2^k$ , k prime or power of prime, characterize the size scales of the space-time sheets in the hierarchy. The p-adic length scale L(k) can be expressed in terms of cell membrane thickness as

$$L(k) = 2^{(k-151)/2} \times L(151) \quad , \tag{1}$$

 $L(151) \simeq 10$  nm. These are so called primary p-adic length scales but there are also n-ary padic length scales related by a scaling of power of  $\sqrt{p}$  to the primary p-adic length scale. Quite recent model for photosynthesis [K6] gives additional support for the importance of also n-ary p-adic length scales so that the relevant p-adic length scales would come as half-octaves in a good approximation but prime and power of prime values of k would be especially important.

### 2.3 Topological light rays (massless extremals, MEs)

I have described MEs, or "topological light rays", in detail in [32] and in [J4, J7], and describe here only very briefly the basic characteristics of MEs and concentrate on new idea about their possible role for consciousness and life.

#### 2.3.1 What MEs are?

MEs (massless extremals, topological light rays) can be regarded as topological field quanta of classical radiation fields [J4, J7]. They are typically tubular space-time sheets inside which radiation fields propagate with light velocity in single direction without dispersion. The simplest case corresponds to a straight cylindrical ME but also curved MEs, kind of curved light rays, are possible. The initial values for a given moment of time are arbitrary by light likeness. Therefore MEs are ideal for precisely targeted communications. What distinguishes MEs from Maxwellian radiation fields in empty space is that light like vacuum 4-current is possible: ordinary Maxwell's equations would state that this current vanishes. Quite generally, purely geometric vacuum charge densities and 3-currents are purely TGD based prediction and could be seen as a classical correlate of the vacuum polarization predicted by quantum field theories.

MEs are fractal structures containing MEs within MEs. The so called scaling law of homeopathy predicts that the high frequency MEs inside low frequency MEs are in a ratio having discrete values [K5]. One can indeed justify this relationship. As ions drop from smaller space-time sheets to magnetic flux tubes, zero point kinetic energy is liberated as high frequency MEs, and the ions dropped to magnetic flux tubes generate cyclotron radiation, and the ratio of the fundamental frequencies is constant not depending on particle mass and being determined solely by p-adic length scale hypothesis. The model for the radio waves induced by the irradiation of DNA by laser light [29] gives support for this picture [K2].

#### 2.3.2 Two basic types of MEs

MEs have 2-dimensional  $CP_2$  projection which means that electro-weak holonomy group is Abelian (color holonomy is always Abelian which suggests that physical states in TGD Universe correspond to states of color multiplets with vanishing color hypercharge and isospin rather than color singlets). If  $CP_2$  projection belongs to a homologically non-trivial geodesic sphere, only em and  $Z^0$  fields and Abelian color gauge fields are present. In the homologically trivial case only classical W fields are non-vanishing.

- 1. Neutral MEs can be assigned to various kinds of communications from biological body to the magnetic body and fractal hierarchy of EEGs and ZEGs represent the basic example in this respect [M3].
- 2. Dark W MEs serving as correlate for dark W exchanges induce an exotic ionization of atomic nuclei [F8, F9, M3]. This induces charge entanglement between magnetic body and biological body generating dark plasma oscillation patterns inducing nerve pulse patterns and ion waves at the space-time sheets occupied by the ordinary matter. The mechanism is based on many-sheeted Faraday law inducing electromagnetic fields at ordinary space-time sheet in turn giving rise to ohmic currents. State function reduction selects one of the exotically

ionized configurations. This mechanism is the most plausible candidate for how magnetic body as an intentional agent controls biological body.

#### 2.3.3 Negative energy MEs

MEs can have either positive or negative energy depending on the time orientation. The understanding of negative energy MEs has increased considerably. Phase conjugate laser beams [24] are the most plausible standard physics counterparts of negative energy MEs since they can be interpreted as time reversed laser beams and do not possess direct Maxwellian analog. By quantumclassical correspondence one can interpret the frequencies associated with negative energy MEs as energies. One can also assume that the Bose-Einstein condensed photons associated with negative energy MEs and with the coherent light generated by the light like vacuum current have negative energies.

For frequencies for which energy is above the thermal energy there is no system which could interact with negative energy MEs or absorb negative energy photons. Therefore negative energy MEs and corresponding photons should propagate through matter practically without any interaction. Feinberg has demonstrated that phase conjugate laser beams behave similarly: for instance, one can see through chickens using these laser beams [25]. This means that negative energy MEs do not respect Faraday cages and thus represent an attractive candidate for the hypothetical Psi field.

Negative energy MEs have many applications.

- 1. Negative energy MEs ideal for generating time like entanglement. Since negative energies are involved, this entanglement can be seen as a correlate for the bound state entanglement leading to a macro-temporal quantum coherence. Negative energy MEs make thus possible telepathic sharing of mental images. Negative energy MEs are involved with both sensory perception, long term memory, and motor action. In the model for living matter [M3] The charge entanglement generated by W MEs inducing exotic weak charge and electromagnetic charge is assumed to be responsible for bio-control whereas neutral MEs in general carrying both em and  $Z^0$  fields are responsible for communications.
- 2. Negative energy MEs are ideal for a precisely targeted realization of intentions. p-Adic ME having a large number of common rational points with negative energy ME is generated and transformed to a real ME in quantum jump. The system receives positive energy and momentum as a recoil effect and the transition is not masked by ordinary spontaneously occurring quantum transitions since the energy of the system increases. One can say that negative energy ME represents the desires communicated to the geometric past and inducing as a reaction the desired action realized as say neuronal activity and generation of positive energy MEs.
- 3. The generation of negative energy MEs is also in a key role in remote metabolism and MEs serve as quantum credit cards implying an extreme flexibility of the metabolism. If the system receiving negative energy MEs is a population inverted laser or its many-sheeted counterpart, then quite a small field intensity associated with negative energy MEs (intensity of negative energy photons) can lead to the amplification of the time reflected positive energy signal. The reason is that the rate for the induced emission is proportional to the number of particles dropped to the ground state from the excited state. Therefore even negative energy bio-photons might serve as quantum controllers of metabolism and induce much more intense beams of positive energy photons, say when interacting with mitochondria.

#### 2.4 Magnetic flux quanta and electrets

Magnetic flux tubes and electrets are extremals of Kähler action dual to each other. Also layer like magnetic flux quanta and their electric counterparts are possible. The magnetic/electric field is in a good approximation of constant magnitude but has varying direction.

#### 2.4.1 Magnetic fields and life

The magnetic field associated with any material system is topologically quantized, and one can assign to any system a magnetic body. An attractive idea is that the relationship of the magnetic body to the material system is to some degree that of the manual to an electronic instrument. Quantitative arguments related to the dark matter hierarchy assuming that magnetic bodies are dark suggest that cognitions and emotions are regarded as somatosensory qualia of the magnetic body [M3, K3]. Magnetic body would in this case serve as a kind of computer screen at which the data items processes in say brain are communicated either classically (positive energy MEs) or by sharing of mental images (negative energy MEs).

Magnetic body is also an active intentional agent: motor actions are controlled from magnetic body and proceed as cascade like processes from long to short length and time scales as quantum communications of desires at various levels of hierarchy of magnetic bodies. Communication occurs backwards in geometric time by negative energy MEs. Motor action as a response to these desires occurs by classical communications by positive energy MEs and as neural activities. This explains the coherence and synchrony of motor actions difficult to understand in neuroscience framework. The sizes of flux quanta are astrophysical: for instance, EEG frequency of 7.8 Hz corresponds to a wave length defined by Earth's circumference. The non-locality in the length scale of magnetosphere, and even in length scales up to light life, is forced by Uncertainty Principle alone, if taken seriously in macroscopic length scales.

The leakage of supra currents of ions and their Cooper pairs from magnetic flux tubes of the Earth's magnetic field to smaller space-time sheets and their dropping back involving liberation of the zero point kinetic energy defines one particular metabolic "Karma's cycle". The dropping of protons from k = 137 atomic space-time sheet involved with the utilization of ATP molecules is only a special instance of the general mechanism involving an entire hierarchy of zero point kinetic energies defining universal metabolic currencies. This leads to the idea that the topologically quantized magnetic field of Earth defines the analog of central nervous system and blood circulation present already during the pre-biotic evolution and making possible primitive metabolism. This has far reaching implications for the understanding of how pre-biotic evolution led to living matter as we understand it [N4].

For instance, it has recently become clear that the dropping of atoms and molecules from space-time sheets labelled by p-adic prime  $p \simeq 2^k$ , k = 131, liberates photons at visible and near infrared wave lengths. The hot k = 131 space-time sheets (with temperatures above 1000 K) could have served as a source of metabolic energy for life-forms at cool k = 137 sheets. Photosynthesis could have developed in the circumstances where solar radiation was replaced with these photons. The correct prediction is that chlorophylls should be especially sensitive to these wave lengths. In particular, it is predicted that also IR wave lengths 700-1000 nm should have been utilized. There indeed are bacteria using only this portion of solar radiation. This leads to a scenario making sense only in TGD universe. Pre-biotic life could have developed at the cool space-time sheets in the hot interior of Earth below crust, where k = 131 space-time sheets are possible and this life could still be there [N4]. Also the life as we know it, could involve hot spots generated by the cavitation of water inside cell. The classical repulsive  $Z^0$  force causes a strong acceleration during final stages of bubble collapse creating high temperatures, and could explain also sono-luminescence [27] as suggested in [F9].

Magnetic Mother Gaia could also form sensory and other representations receiving input from

several brains via negative energy EEG MEs entangling magnetosphere with brains. The multibrained magnetospheric selves could be responsible for the third person aspect of consciousness and for the evolution of social structures. For instance, the successful healing by prayer and meditation groups [33], and the experiments of Mark Germine [34] provide support for the notion of multibrained magnetospheric selves are involved. Magnetic flux tubes could function as wave guides for MEs and this aspect is crucial in the model of long term memory.

#### 2.4.2 Electrets and bio-systems

Bio-systems are known to be full of electrets and liquid crystals [35]. Perhaps the most fundamental electret structure is cell membrane. In particular, the water inside cells tends to be in gel phase which is liquid crystal phase. There are many good reasons for why water should be in ordered phase. One very fundamental reason is that bio-polymers are stable in liquid crystal/ordered water phase since there are no free water molecules available for the depolymerization by hydration. In fact, only a couple of years ago it was experimentally discovered that bio-polymers can be stabilized around ice.

The capacitor model for sensory receptor is one very important application of the electret concept [31, K3]. Sensory qualia result in the flow of particles with given quantum numbers from the plate to another one in quantum discharge. This kind of amplification of quantum number *resp.* zero mode increments would give rise to both geometric *resp.* non-geometric qualia [K3].

Also micro-tubuli are electrets. Sol-gel transition, as any phase transition, is an good candidate for the representation of a conscious bit and controlled local sol-gel transitions between ordinary and liquid crystal water could be a basic control tool making possible cellular locomotion, changes of protein conformations, etc... The tubulin dimers of micro-tubuli could induce sol-gel transformations by generating negative energy MEs, and micro-tubular surface could provide bit maps of their environment somewhat like sensory areas of brain provide maps of body. If gel $\rightarrow$ sol transition around tubulin inducing conformational change induces sol $\rightarrow$ gel transformation in some point of environment as would be the case for the seesaw mechanism to be discussed below, a one-one correspondence would result. By this one-one correspondence micro-tubules would automatically generate kind of conscious log files about the control activities which could have evolved to micro-tubular declarative memory representations about what happens inside cell [K6].

# 3 General considerations

In this section field equations and their physical interpretation are discussed. Quantum classical correspondence suggests that the non-deterministic dynamics of Kähler action makes possible self-referential dynamics in the sense that larger space-time sheet perform smoothed out mimicry of the dynamics at smaller space-time sheets. The fact that the divergence of the energy momentum tensor, Lorentz 4-force, does not vanish in general makes possible the mimicry of even dissipation and of the second law. For asymptotic self organization patterns for which dissipation is absent the Lorentz 4-force must vanish. This condition is guaranteed if Kähler current is proportional to the instanton current in the case that  $CP_2$  projection of the space-time sheet is smaller than four and vanishes otherwise. An attractive identification for the vanishing of Lorentz 4-force is as a condition equivalent with the principle selecting preferred extremals of Kähler action so that this principle would be essentially equivalent with the second law of thermodynamics.

# 3.1 Long range classical weak and color gauge fields as correlates for dark massless weak bosons

Long ranged classical electro-weak gauge fields are unavoidably present when the dimension D of the  $CP_2$  projection of the space-time sheet is larger than 2. Classical color gauge fields are non-vanishing for all non-vacuum extremals. This poses deep interpretational problems. If ordinary quarks and leptons are assumed to carry weak charges feeded to larger space-time sheets within electro-weak length scale, large hadronic, nuclear, and atomic parity breaking effects, large contributions of the classical  $Z^0$  force to Rutherford scattering, and strong isotopic effects, are expected. If weak charges are screened within electro-weak length scale, the question about the interpretation of long ranged classical weak fields remains.

#### 3.1.1 Various interpretations for the long ranged classical electro-weak fields

During years I have discussed several solutions to the problems listed above.

Option I: The trivial solution of the constraints is that  $Z^0$  charges are neutralized at electroweak length scale. The problem is that this option leaves open the interpretation of classical long ranged electro-weak gauge fields unavoidably present in all length scales when the dimension for the  $CP_2$  projection of the space-time surface satisfies D > 2.

Option II: Second option involves several variants but the basic assumption is that nuclei or even quarks feed their  $Z^0$  charges to a space-time sheet with size of order neutrino Compton length. The large parity breaking effects in hadronic, atomic, and nuclear length scales is not the only difficulty. The scattering of electrons, neutrons and protons in the classical long range  $Z^0$  force contributes to the Rutherford cross section and it is very difficult to see how neutrino screening could make these effects small enough. Strong isotopic effects in condensed matter due to the classical  $Z^0$  interaction energy are expected. It is far from clear whether all these constraints can be satisfied by any assumptions about the structure of topological condensate.

*Option III*: During 2005 (27 years after the birth of TGD!) third option solving the problems emerged based on the progress in the understanding of the basic mathematics behind TGD.

In ordinary phase the  $Z^0$  charges of elementary particles are indeed neutralized in intermediate boson length scale so that the problems related to the parity breaking, the large contributions of classical  $Z^0$  force to Rutherford scattering, and large isotopic effects in condensed matter, trivialize.

Classical electro-weak gauge fields in macroscopic length scales are identified as space-time correlates for the gauge fields created by dark matter, which corresponds to a macroscopically quantum coherent phase for which elementary particles possess complex conformal weights such that the net conformal weight of the system is real.

In this phase  $U(2)_{ew}$  symmetry is not broken below the scaled up weak scale except for fermions so that gauge bosons are massless below this length scale whereas fermion masses are essentially the same as for ordinary matter. By charge screening gauge bosons look massive in length scales much longer than the relevant p-adic length scale. The large parity breaking effects in living matter (chirality selection for bio-molecules) support the view that dark matter is what makes living matter living.

#### 3.1.2 Classical color gauge fields

Classical long ranged color gauge fields always present for non-vacuum extremals are interpreted as space-time correlates of gluon fields associated with dark copies of hadron physics. It seems that this picture is indeed what TGD predicts. A very special feature of classical color fields is that the holonomy group is Abelian. This follows directly from the expression  $g^A_{\alpha\beta} \propto H^A J_{\alpha\beta}$  of induced gluon fields in terms of Hamiltonians  $H^A$  of color isometries and induced Kähler form  $J_{\alpha\beta}$ . This means that classical color magnetic and electric fluxes reduce to the analogs of ordinary magnetic fluxes appearing in the construction of configuration space geometry [B2, B3].

By a local color rotation the color field can be rotated to a fixed direction so that genuinely Abelian field would be in question apart from the possible presence of gauge singularities making impossible a global selection of color direction. These singularities could be present since Kähler form defines a magnetic monopole field. An interesting question inspired by quantum classical correspondence is what the Abelian holonomy tells about the sources of color gauge fields and color confinement.

For instance, could Abelian color holonomy mean that colored gluons (and presumably also other colored particles) do not propagate in the p-adic length scale considered? Color neutral gluons would propagate but since also their sources must be color neutral, they should have vanishing net color electric fluxes. This form of confinement would allow those states of color multiplets which have vanishing color charges and obviously symmetry breaking down to  $U(1) \times U(1)$  would be in question. This would serve as a signal for monopole confinement which would not exclude higher multipoles for the Abelian color fields. This kind of fields appear in the TGD based model for nuclei as nuclear strings bound together by color flux tubes [F8].

### 3.2 Is absolute minimization the correct variational principle?

One can criticize the original assumption that extremals correspond to absolute minima, and the number theoretical vision discussed in [E2] indeed favors the separate minimization of magnitudes of positive and negative contributions to the Kähler action.

For this option Universe would do its best to save energy, being as near as possible to vacuum. Also vacuum extremals would become physically relevant: note that they would be only inertial vacua and carry non-vanishing density gravitational energy. The non-determinism of the vacuum extremals would have an interpretation in terms of the ability of Universe to engineer itself.

The 3-surfaces for which  $CP_2$  projection is at least 2-dimensional and not a Lagrange manifold would correspond to non-vacua since conservation laws do not leave any other option. The variational principle would favor equally magnetic and electric configurations whereas absolute minimization of action based on  $S_K$  would favor electric configurations. The positive and negative contributions would be minimized for 4-surfaces in relative homology class since the boundary of  $X^4$  defined by the intersections with 7-D light-like causal determinants would be fixed. Without this constraint only vacuum bubbles would result.

The attractiveness of the number theoretical variational principle from the point of calculability of TGD would be that the initial values for the time derivatives of the imbedding space coordinates at  $X^3$  at light-like 7-D causal determinant could be computed by requiring that the energy of the solution is minimized. This could mean a computerizable solution to the construction of Kähler function.

Since the considerations of this chapter relate only to the extremals of Kähler action, I have not bothered to replace "absolute minimum" with "preferred extremal" in the text. The number theoretic approach based on the properties of quaternions and octonions discussed in the chapter [E2] leads to a proposal for the general solution of field equations based on the generalization of the notion of calibration [16] providing absolute minima of volume to that of Kähler calibration. This approach will not be discussed in this chapter.

#### 3.3 Field equations

The requirement that Kähler action is stationary leads to the following field equations in the interior of the four-surface

$$D_{\beta}(T^{\alpha\beta}h_{\alpha}^{k}) - j^{\alpha}J_{l}^{k}\partial_{\alpha}h^{l} = 0 ,$$
  

$$T^{\alpha\beta} = J^{\nu\alpha}J_{\nu}^{\ \beta} - \frac{1}{4}g^{\alpha\beta}J^{\mu\nu}J_{\mu\nu} .$$
(2)

Here  $T^{\alpha\beta}$  denotes the traceless canonical energy momentum tensor associated with the Kähler action.

An equivalent form for the first equation is

$$T^{\alpha\beta}H^{k}_{\alpha\beta} - j^{\alpha}(J^{\beta}_{\alpha}h^{k}_{\beta} + J^{k}_{l}\partial_{\alpha}h^{l}) = 0 .$$
  

$$H^{k}_{\alpha\beta} = D_{\beta}\partial_{\alpha}h^{k} .$$
(3)

 $H^k_{\alpha\beta}$  denotes the components of the second fundamental form and  $j^{\alpha} = D_{\beta}J^{\alpha\beta}$  is the gauge current associated with the Kähler field.

On the boundaries of  $X^4$  the field equations are given by the expression

$$T^{n\beta}\partial_{\beta}h^{k} - J^{n\alpha}(J^{\ \beta}_{\alpha}\partial_{\beta}h^{k} + J^{k}_{\ l})\partial_{\alpha}h^{k}) = 0 \quad . \tag{4}$$

A general manner manner to solve the field equations on the boundaries is to assume that the induced Kähler field associated with the boundaries vanishes:

$$J_{\alpha\beta}(\delta) = 0 . (5)$$

In this case the energy-momentum tensor vanishes identically on the boundary component. On the outer boundaries of the 3-surface this solution ansatz makes sense only provided the gauge fluxes and gravitational flux (defined by Newtonian potential in the non-relativistic limit) associated with the matter in the interior go somewhere. The only possibility seems to be that 3-surface is topologically condensed on a larger 3-surface and feeds its gauge fluxes to the larger 3-surface via # contacts (topological sum). This assumption forces the concept of topological condensate defined as a hierarchical structure of 3-surfaces condensed on each other and thus giving rise to the many-sheeted space-time.

An important thing to notice is that the boundary conditions do not force the normal components of the gauge fields to zero even if the Kähler electric field vanishes near the boundaries. This makes in principle possible gauge charge renormalization classically resulting from the hierarchical structure of the topological condensation.

## 3.4 Topologization and light-likeness of the Kähler current as alternative manners to guarantee vanishing of Lorentz 4-force

The general solution of 4-dimensional Einstein-Yang Mills equations in Euclidian 4-metric relies on self-duality of the gauge field, which topologizes gauge charge. This topologization can be achieved by a weaker condition, which can be regarded as a dynamical generalization of the Beltrami condition. An alternative manner to achieve vanishing of the Lorentz 4-force is light-likeness of the Kähler 4-current. This does not require topologization.

#### **3.4.1** Topologization of the Kähler current for $D_{CP_2} = 3$ : covariant formulation

The condition states that Kähler 4-current is proportional to the instanton current whose divergence is instanton density and vanishes when the dimension of  $CP_2$  projection is smaller than four:  $D_{CP_2} < 4$ . For  $D_{CP_2} = 2$  the instanton 4-current vanishes identically and topologization is equivalent with the vanishing of the Kähler current.

$$j^{\alpha} \equiv D_{\beta} J^{\alpha\beta} = \psi \times j^{\alpha}_{I} = \psi \times \epsilon^{\alpha\beta\gamma\delta} J_{\beta\gamma} A_{\delta} \quad . \tag{6}$$

Here the function  $\psi$  is an arbitrary function  $\psi(s^k)$  of  $CP_2$  coordinates  $s^k$  regarded as functions of space-time coordinates. It is essential that  $\psi$  depends on the space-time coordinates through the  $CP_2$  coordinates only. Hence the representation as an imbedded gauge field is crucial element of the solution ansatz.

The field equations state the vanishing of the divergence of the 4-current. This is trivially true for instanton current for  $D_{CP_2} < 4$ . Also the contraction of  $\nabla \psi$  (depending on space-time coordinates through  $CP_2$  coordinates only) with the instanton current is proportional to the winding number density and therefore vanishes for  $D_{CP_2} < 4$ .

The topologization of the Kähler current guarantees the vanishing of the Lorentz 4-force. Indeed, using the self-duality condition for the current, the expression for the Lorentz 4-force reduces to a term proportional to the instanton density:

$$j^{\alpha}J_{\alpha\beta} = \psi \times j_{I}^{\alpha}J_{\alpha\beta}$$
$$= \psi \times \epsilon^{\alpha\mu\nu\delta}J_{\mu\nu}A_{\delta}J_{\alpha\beta} .$$
(7)

Since all vector quantities appearing in the contraction with the four-dimensional permutation tensor are proportional to the gradients of  $CP_2$  coordinates, the expression is proportional to the instanton density, and thus winding number density, and vanishes for  $D_{CP_2} < 4$ .

Remarkably, the topologization of the Kähler current guarantees also the vanishing of the term  $j^{\alpha}J^{k_l}\partial_{\alpha}s^k$  in the field equations for  $CP_2$  coordinates. This means that field equations reduce in both  $M_+^4$  and  $CP_2$  degrees of freedom to

$$T^{\alpha\beta}H^k_{\alpha\beta} = 0 . aga{8}$$

These equations differ from the equations of minimal surface only by the replacement of the metric tensor with energy momentum tensor. The earlier proposal that quaternion conformal invariance in a suitable sense might provide a general solution of the field equations could be seen as a generalization of the ordinary conformal invariance of string models. If the topologization of the Kähler current implying effective dimensional reduction in  $CP_2$  degrees of freedom is consistent with quaternion conformal invariance, the quaternion conformal structures must differ for the different dimensions of  $CP_2$  projection.

#### **3.4.2** Topologization of the Kähler current for $D_{CP_2} = 3$ : non-covariant formulation

In order to gain a concrete understanding about what is involved it is useful to repeat these arguments using the 3-dimensional notation. The components of the instanton 4-current read in three-dimensional notation as

$$\overline{j}_I = \overline{E} \times \overline{A} + \phi \overline{B} , \quad \rho_I = \overline{B} \cdot \overline{A} .$$
 (9)

The self duality conditions for the current can be written explicitly using 3-dimensional notation and read

$$\nabla \times \overline{B} - \partial_t \overline{E} = \overline{j} = \psi \overline{j}_I = \psi \left( \phi \overline{B} + \overline{E} \times \overline{A} \right) ,$$
  

$$\nabla \cdot E = \rho = \psi \rho_I .$$
(10)

For a vanishing electric field the self-duality condition for Kähler current reduces to the Beltrami condition

$$\nabla \times \overline{B} = \alpha \overline{B} , \quad \alpha = \psi \phi .$$
 (11)

The vanishing of the divergence of the magnetic field implies that  $\alpha$  is constant along the field lines of the flow. When  $\phi$  is constant and  $\overline{A}$  is time independent, the condition reduces to the Beltrami condition with  $\alpha = \phi = constant$ , which allows an explicit solution [19].

One can check also the vanishing of the Lorentz 4-force by using 3-dimensional notation. Lorentz 3-force can be written as

$$\rho_I \overline{E} + \overline{j} \times \overline{B} = \psi \overline{B} \cdot \overline{AE} + \psi \left( \overline{E} \times \overline{A} + \phi \overline{B} \right) \times \overline{B} = 0 \quad . \tag{12}$$

The fourth component of the Lorentz force reads as

$$\overline{j} \cdot \overline{E} = \psi \overline{B} \cdot \overline{E} + \psi \left( \overline{E} \times \overline{A} + \phi \overline{B} \right) \cdot \overline{E} = 0 \quad . \tag{13}$$

The remaining conditions come from the induction law of Faraday and could be guaranteed by expressing  $\overline{E}$  and  $\overline{B}$  in terms of scalar and vector potentials.

The density of the Kähler electric charge of the vacuum is proportional to the helicity density of the so called helicity charge  $\rho = \psi \rho_I = \psi B \cdot A$ . This charge is topological charge in the sense that it does not depend on the induced metric at all. Note the presence of arbitrary function  $\psi$  of  $CP_2$  coordinates.

Further conditions on the functions appearing in the solution ansatz come from the 3 independent field equations for  $CP_2$  coordinates. What is remarkable that the generalized self-duality condition for the Kähler current allows to understand the general features of the solution ansatz to very high degree without any detailed knowledge about the detailed solution. The question whether field equations allow solutions consistent with the self duality conditions of the current will be dealt later. The optimistic guess is that the field equations and topologization of the Kähler current relate to each other very intimately.

# 3.4.3 Vanishing or light likeness of the Kähler current guarantees vanishing of the Lorentz 4-force for $D_{CP_2} = 2$

For  $D_{CP_2} = 2$  one can always take two  $CP_2$  coordinates as space-time coordinates and from this it is clear that instanton current vanishes so that topologization gives a vanishing Kähler current. In particular, the Beltrami condition  $\nabla \times \overline{B} = \alpha \overline{B}$  is not consistent with the topologization of the instanton current for  $D_{CP_2} = 2$ .

 $D_{CP_2} = 2$  case can be treated in a coordinate invariant manner by using the two coordinates of  $CP_2$  projection as space-time coordinates so that only a magnetic or electric field is present depending on whether the gauge current is time-like or space-like. Light-likeness of the gauge current provides a second manner to achieve the vanishing of the Lorentz force and is realized in case of massless extremals having  $D_{CP_2} = 2$ : this current is in the direction of propagation whereas magnetic and electric fields are orthogonal to it so that Beltrami conditions is certainly not satisfied.

#### 3.4.4 Under what conditions topologization of Kähler current yields Beltrami conditions?

Topologization of the Kähler 4-current gives rise to magnetic Beltrami fields if either of the following conditions is satisfied.

1. The  $\overline{E} \times \overline{A}$  term contributing besides  $\phi \overline{B}$  term to the topological current vanishes. This requires that  $\overline{E}$  and  $\overline{A}$  are parallel to each other

$$\overline{E} = \nabla \Phi - \partial_t \overline{A} = \beta \overline{A} \tag{14}$$

This condition is analogous to the Beltrami condition. Now only the 3-space has as its coordinates time coordinate and two spatial coordinates and and  $\overline{B}$  is replaced with  $\overline{A}$ . Since E and B are orthogonal, this condition implies  $\overline{B} \cdot \overline{A} = 0$  so that Kähler charge density is vanishing.

2. The vector  $\overline{E} \times \overline{A}$  is parallel to  $\overline{B}$ .

$$\overline{E} \times \overline{A} = \beta \overline{B} \tag{15}$$

The condition is consistent with the orthogonality of  $\overline{E}$  and  $\overline{B}$  but implies the orthogonality of  $\overline{A}$  and  $\overline{B}$  so that electric charge density vanishes.

In both cases vector potential fails to define a contact structure since  $B \cdot A$  vanishes (contact structures are discussed briefly below), and there exists a global coordinate along the field lines of  $\overline{A}$  and the full contact structure is lost again. Note however that the Beltrami condition for magnetic field means that magnetic field defines a contact structure irrespective of whether  $\overline{B} \cdot \overline{A}$  vanishes or not. The transition from the general case to Beltrami field would thus involve the replacement

$$(\overline{A},\overline{B}) \to_{\nabla \times} (\overline{B},\overline{j})$$

induced by the rotor.

One must of course take these considerations somewhat cautiously since the inner product depends on the induced 4-metric and it might be that induced metric could allow small vacuum charge density and make possible genuine contact structure.

#### 3.4.5 Hydrodynamic analogy

The field equations of TGD are basically hydrodynamic equations stating the local conservation of the currents associated with the isometries of the imbedding space. Therefore it is intriguing that Beltrami fields appear also as solutions of ideal magnetohydrodynamics equations and as steady solutions of non-viscous incompressible flow described by Euler equations [20].

In hydrodynamics the role of the magnetic field is taken by the velocity field. TGD based models for nuclei [F6] and condensed matter [F9] involve in an essential manner valence quarks having large  $\hbar$  and exotic quarks giving nucleons anomalous color and weak charges creating long ranged color and weak forces. Weak forces have a range of order atomic radius and could be responsible for the repulsive core in van der Waals potential.

This raises the idea that the incompressible flow could occur along the field lines of the  $Z^0$  magnetic field so that the velocity field would be proportional to the  $Z^0$  magnetic field and the

Beltrami condition for the velocity field would reduce to that for  $Z^0$  magnetic field. Thus the flow lines of hydrodynamic flow would directly correspond to those of  $Z^0$  magnetic field. The generalized Beltrami flow based on the topologization of the  $Z^0$  current would allow to model also the non-stationary incompressible non-viscous hydrodynamical flows.

It would seem that one cannot describe viscous flows using flows satisfying generalized Beltrami conditions since the vanishing of the Lorentz 4-force says that there is no local dissipation of the classical field energy. One might claim that this is not a problem since in TGD framework viscous flow could be seen as a practical description of a quantum jump sequence by replacing the corresponding sequence of space-time surfaces with a single space-time surface.

One the other hand, quantum classical correspondence requires that also dissipative effects have space-time correlates. Kähler fields, which are dissipative, and thus correspond to a nonvanishing Lorentz 4-force, represent one candidate for correlates of this kind. If this is the case, then the fields satisfying the generalized Beltrami condition provide space-time correlates only for the asymptotic self organization patterns for which the viscous effects are negligible, and also the solutions of field equations describing effects of viscosity should be possible.

One must however take this argument with a grain of salt. Dissipation, that is the transfer of conserved quantities to degrees of freedom corresponding to shorter scales, could correspond to a transfer of these quantities between different space-time sheets of the many-sheeted spacetime. Here the opponent could however argue that larger space-time sheets mimic the dissipative dynamics in shorter scales and that classical currents represent "symbolically" averaged currents in shorter length scales, and that the local non-conservation of energy momentum tensor consistent with local conservation of isometry currents provides a unique manner to mimic the dissipative dynamics. This view will be developed in more detail below.

#### 3.4.6 The stability of generalized Beltrami fields

The stability of generalized Beltrami fields is of high interest since unstable points of space-time sheets are those around which macroscopic changes induced by quantum jumps are expected to be localized.

#### 1. Contact forms and contact structures

The stability of Beltrami flows has been studied using the theory of contact forms in threedimensional Riemann manifolds [21]. Contact form is a one-form A (that is covariant vector field  $A_{\alpha}$ ) with the property  $A \wedge dA \neq 0$ . In the recent case the induced Kähler gauge potential  $A_{\alpha}$  and corresponding induced Kähler form  $J_{\alpha\beta}$  for any 3-sub-manifold of space-time surface define a contact form so that the vector field  $A^{\alpha} = g^{\alpha\beta}A_{\beta}$  is not orthogonal with the magnetic field  $B^{\alpha} = \epsilon^{\alpha\beta\delta}J_{\beta\gamma}$ . This requires that magnetic field has a helical structure. Induced metric in turn defines the Riemann structure.

If the vector potential defines a contact form, the charge density associated with the topologized Kähler current must be non-vanishing. This can be seen as follows.

- 1. The requirement that the flow lines of a one-form  $X_{\mu}$  defined by the vector field  $X^{\mu}$  as its dual allows to define a global coordinate x varying along the flow lines implies that there is an integrating factor  $\phi$  such that  $\phi X = dx$  and therefore  $d(\phi X) = 0$ . This implies  $dlog(\phi) \wedge X = -dX$ . From this the necessary condition for the existence of the coordinate xis  $X \wedge dX = 0$ . In the three-dimensional case this gives  $\overline{X} \cdot (\nabla \times \overline{X}) = 0$ .
- 2. This condition is by definition not satisfied by the vector potential defining a contact form so that one cannot identify a global coordinate varying along the flow lines of the vector potential. The condition  $\overline{B} \cdot \overline{A} \neq 0$  states that the charge density for the topologized Kähler current is non-vanishing. The condition that the field lines of the magnetic field allow a

global coordinate requires  $\overline{B} \cdot \nabla \times \overline{B} = 0$ . The condition is not satisfied by Beltrami fields with  $\alpha \neq 0$ . Note that in this case magnetic field defines a contact structure.

Contact structure requires the existence of a vector  $\xi$  satisfying the condition  $A(\xi) = 0$ . The vector field  $\xi$  defines a plane field, which is orthogonal to the vector field  $A^{\alpha}$ . Reeb field in turn is a vector field for which A(X) = 1 and dA(X;) = 0 hold true. The latter condition states the vanishing of the cross product  $X \times B$  so that X is parallel to the Kähler magnetic field  $B^{\alpha}$  and has unit projection in the direction of the vector field  $A^{\alpha}$ . Any Beltrami field defines a Reeb field irrespective of the Riemannian structure.

#### 2. Stability of the Beltrami flow and contact structures

Contact structures are used in the study of the topology and stability of the hydrodynamical flows [21], and one might expect that the notion of contact structure and its proper generalization to the four-dimensional context could be useful in TGD framework also. An example giving some idea about the complexity of the flows defined by Beltrami fields is the Beltrami field in  $R^3$  possessing closed orbits with all possible knot and link types simultaneously [21]!

Beltrami flows associated with Euler equations are known to be unstable [21]. Since the flow is volume preserving, the stationary points of the Beltrami flow are saddle points at which also vorticity vanishes and linear instabilities of Navier-Stokes equations can develop. From the point of view of biology it is interesting that the flow is stabilized by vorticity which implies also helical structures. The stationary points of the Beltrami flow correspond in TGD framework to points at which the induced Kähler magnetic field vanishes. They can be unstable by the vacuum degeneracy of Kähler action implying classical non-determinism. For generalized Beltrami fields velocity and vorticity (both divergence free) are replaced by Kähler current and instanton current.

More generally, the points at which the Kähler 4-current vanishes are expected to represent potential instabilities. The instanton current is linear in Kähler field and can vanish in a gauge invariant manner only if the induced Kähler field vanishes so that the instability would be due to the vacuum degeneracy also now. Note that the vanishing of the Kähler current allows also the generation of region with  $D_{CP_2} = 4$ . The instability of the points at which induce Kähler field vanish is manifested in quantum jumps replacing the generalized Beltrami field with a new one such that something new is generated around unstable points. Thus the regions in which induced Kähler field becomes weak are the most interesting ones. For example, unwinding of DNA could be initiated by an instability of this kind.

#### 3.5 How to satisfy field equations?

The topologization of the Kähler current guarantees also the vanishing of the term  $j^{\alpha}J^{k_l}\partial_{\alpha}s^k$  in the field equations for  $CP_2$  coordinates. This means that field equations reduce in both  $M_+^4$  and  $CP_2$  degrees of freedom to

$$T^{\alpha\beta}H^k_{\alpha\beta} = 0 . (16)$$

These equations differ from the equations of minimal surface only by the replacement of the metric tensor with energy momentum tensor. The earlier proposal that quaternion conformal invariance in a suitable sense might provide a general solution of the field equations could be seen as a generalization of the ordinary conformal invariance of string models. If the topologization of the Kähler current implying effective dimensional reduction in  $CP_2$  degrees of freedom is consistent with quaternion conformal invariance, the quaternion conformal structures must differ for the different dimensions of  $CP_2$  projection. In the following somewhat different approach is however considered utilizing the properties of Hamilton Jacobi structures of  $M_+^4$  introduced in the study of massless extremals and contact structures of  $CP_2$  emerging naturally in the case of generalized Beltrami fields.

#### 3.5.1 String model as a starting point

String model serves as a starting point.

- 1. In the case of Minkowskian minimal surfaces representing string orbit the field equations reduce to purely algebraic conditions in light cone coordinates (u, v) since the induced metric has only the component  $g_{uv}$ , whereas the second fundamental form has only diagonal components  $H_{uu}^k$  and  $H_{vv}^k$ .
- 2. For Euclidian minimal surfaces (u, v) is replaced by complex coordinates  $(w, \overline{w})$  and field equations are satisfied because the metric has only the component  $g^{w\overline{w}}$  and second fundamental form has only components of type  $H_{ww}^k$  and  $H_{\overline{ww}}^{\overline{k}}$ . The mechanism should generalize to the recent case.

# **3.5.2** The general form of energy momentum tensor as a guideline for the choice of coordinates

Any 3-dimensional Riemann manifold allows always a orthogonal coordinate system for which the metric is diagonal. Any 4-dimensional Riemann manifold in turn allows a coordinate system for which 3-metric is diagonal and the only non-diagonal components of the metric are of form  $g^{ti}$ . This kind of coordinates might be natural also now. When  $\overline{E}$  and  $\overline{B}$  are orthogonal, energy momentum tensor has the form

$$T = \begin{pmatrix} \frac{E^2 + B^2}{2} & 0 & 0 & EB \\ 0 & \frac{E^2 + B^2}{2} & 0 & 0 \\ 0 & 0 & \frac{-E^2 + B^2}{2} & 0 \\ EB & 0 & 0 & \frac{E^2 - B^2}{2} \end{pmatrix}$$
(17)

in the tangent space basis defined by time direction and longitudinal direction  $\overline{E} \times \overline{B}$ , and transversal directions  $\overline{E}$  and  $\overline{B}$ . Note that T is traceless.

The optimistic guess would be that the directions defined by these vectors integrate to three orthogonal coordinates of  $X^4$  and together with time coordinate define a coordinate system containing only  $g^{ti}$  as non-diagonal components of the metric. This however requires that the fields in question allow an integrating factor and, as already found, this requires  $\nabla \times X \cdot X = 0$  and this is not the case in general.

Physical intuition suggests however that  $X^4$  coordinates allow a decomposition into longitudinal and transversal degrees freedom. This would mean the existence of a time coordinate t and longitudinal coordinate z the plane defined by time coordinate and vector  $\overline{E} \times \overline{B}$  such that the coordinates u = t - z and v = t + z are light like coordinates so that the induced metric would have only the component  $g^{uv}$  whereas  $g^{vv}$  and  $g^{uu}$  would vanish in these coordinates. In the transversal space-time directions complex space-time coordinate coordinate w could be introduced. Metric could have also non-diagonal components besides the components  $g^{w\overline{w}}$  and  $g^{uv}$ .

#### **3.5.3** Hamilton Jacobi structures in $M^4_+$

Hamilton Jacobi structure in  $M^4_+$  can understood as a generalized complex structure combing transversal complex structure and longitudinal hyper-complex structure so that notion of holomorphy and Kähler structure generalize.

- 1. Denote by  $m^i$  the linear Minkowski coordinates of  $M^4$ . Let  $(S^+, S^-, E^1, E^2)$  denote local coordinates of  $M^4_+$  defining a *local* decomposition of the tangent space  $M^4$  of  $M^4_+$  into a direct, not necessarily orthogonal, sum  $M^4 = M^2 \oplus E^2$  of spaces  $M^2$  and  $E^2$ . This decomposition has an interpretation in terms of the longitudinal and transversal degrees of freedom defined by local light-like four-velocities  $v_{\pm} = \nabla S_{\pm}$  and polarization vectors  $\epsilon_i = \nabla E^i$  assignable to light ray. Assume that  $E^2$  allows complex coordinates  $w = E^1 + iE^2$  and  $\overline{w} = E^1 iE^2$ . The simplest decomposition of this kind corresponds to the decomposition  $(S^+ \equiv u = t + z, S^- \equiv v = t z, w = x + iy, \overline{w} = x iy)$ .
- 2. In accordance with this physical picture,  $S^+$  and  $S^-$  define light-like curves which are normals to light-like surfaces and thus satisfy the equation:

$$(\nabla S_{\pm})^2 = 0$$

The gradients of  $S_{\pm}$  are obviously analogous to local light like velocity vectors  $v = (1, \overline{v})$  and  $\tilde{v} = (1, -\overline{v})$ . These equations are also obtained in geometric optics from Hamilton Jacobi equation by replacing photon's four-velocity with the gradient  $\nabla S$ : this is consistent with the interpretation of massless extremals as Bohr orbits of em field.  $S_{\pm} = constant$  surfaces can be interpreted as expanding light fronts. The interpretation of  $S_{\pm}$  as Hamilton Jacobi functions justifies the term Hamilton Jacobi structure.

The simplest surfaces of this kind correspond to t = z and t = -z light fronts which are planes. They are dual to each other by hyper complex conjugation  $u = t - z \rightarrow v = t + z$ . One should somehow generalize this conjugation operation. The simplest candidate for the conjugation  $S^+ \rightarrow S^-$  is as a conjugation induced by the conjugation for the arguments:  $S^+(t-z, t+z, x, y) \rightarrow S^-(t-z, t+z, x, y) = S^+(t+z, t-z, x, -y)$  so that a dual pair is mapped to a dual pair. In transversal degrees of freedom complex conjugation would be involved.

3. The coordinates  $(S_{\pm}, w, \overline{w})$  define local light cone coordinates with the line element having the form

$$ds^{2} = g_{+-}dS^{+}dS^{-} + g_{w\overline{w}}dwd\overline{w} + g_{+w}dS^{+}dw + g_{+\overline{w}}dS^{+}d\overline{w} + g_{-w}dS^{-}dw + g_{-\overline{w}}dS^{-}d\overline{w} .$$
(18)

Conformal transformations of  $M^4_+$  leave the general form of this decomposition invariant. Also the transformations which reduces to analytic transformations  $w \to f(w)$  in transversal degrees of freedom and hyper-analytic transformations  $S^+ \to f(S^+), S^- \to f(S^-)$  in longitudinal degrees of freedom preserve this structure.

4. The basic idea is that of generalized Kähler structure meaning that the notion of Kähler function generalizes so that the non-vanishing components of metric are expressible as

$$g_{w\overline{w}} = \partial_w \partial_{\overline{w}} K \quad , \qquad g_{+-} = \partial_{S^+} \partial_{S^-} K \quad ,$$

$$g_{w\pm} = \partial_w \partial_{S^\pm} K \quad , \qquad g_{\overline{w}\pm} = \partial_{\overline{w}} \partial_{S^\pm} K \quad .$$
(19)

for the components of the metric. The expression in terms of Kähler function is coordinate invariant for the same reason as in case of ordinary Kähler metric. In the standard lightcone coordinates the Kähler function is given by

$$K = w_0 \overline{w}_0 + uv$$
,  $w_0 = x + iy$ ,  $u = t - z$ ,  $v = t + z$ . (20)

The Christoffel symbols satisfy the conditions

$${k \atop w \overline{w}} = 0$$
,  ${k \atop +-} = 0$ . (21)

If energy momentum tensor has only the components  $T^{w\overline{w}}$  and  $T^{+-}$ , field equations are satisfied in  $M_{+}^{4}$  degrees of freedom.

5. The Hamilton Jacobi structures related by these transformations can be regarded as being equivalent. Since light-like 3- surface is, as the dynamical evolution defined by the light front, fixed by the 2-surface serving as the light source, these structures should be in one-one correspondence with 2-dimensional surfaces with two surfaces regarded as equivalent if they correspond to different time=constant snapshots of the same light front, or are related by a conformal transformation of  $M_+^4$ . Obviously there should be quite large number of them. Note that the generating two-dimensional surfaces relate also naturally to generalized conformal invariance and corresponding Kac Moody invariance for which deformations defined by the  $M^4$  coordinates as functions of the light-cone coordinates of the light front evolution define Kac Moody algebra, which thus seems to appear naturally also at the level of solutions of field equations.

The task is to find all possible local light cone coordinates defining one-parameter families 2surfaces defined by the condition  $S_i = constant$ , i = + or = -, dual to each other and expanding with light velocity. The basic open questions are whether the generalized Kähler function indeed makes sense and whether the physical intuition about 2-surfaces as light sources parameterizing the set of all possible Hamilton Jacobi structures makes sense.

#### **3.5.4** Contact structure and generalized Kähler structure of CP<sub>2</sub> projection

In the case of 3-dimensional  $CP_2$  projection it is assumed that one can introduce complex coordinates  $(\xi, \overline{\xi})$  and the third coordinate s. These coordinates would correspond to a contact structure in 3-dimensional  $CP_2$  projection defining transversal symplectic and Kähler structures. In these coordinates the transversal parts of the induced  $CP_2$  Kähler form and metric would contain only components of type  $g_{w\overline{w}}$  and  $J_{w\overline{w}}$ . The transversal Kähler field  $J_{w\overline{w}}$  would induce the Kähler magnetic field and the components  $J_{sw}$  and  $J_{s\overline{w}}$  the Kähler electric field.

It must be emphasized that the non-integrability of the contact structure implies that J cannot be parallel to the tangent planes of s = constant surfaces, s cannot be parallel to neither A nor the dual of J, and  $\xi$  cannot vary in the tangent plane defined by J. A further important conclusion is that for the solutions with 3-dimensional  $CP_2$  projection topologized Kähler charge density is necessarily non-vanishing by  $A \wedge J \neq 0$  whereas for the solutions with  $D_{CP_2} = 2$  topologized Kähler current vanishes.

Also the  $CP_2$  projection is assumed to possess a generalized Kähler structure in the sense that all components of the metric except  $s_{ss}$  are derivable from a Kähler function by formulas similar to  $M^4_+$  case.

$$s_{w\overline{w}} = \partial_w \partial_{\overline{w}} K$$
,  $s_{ws} = \partial_w \partial_s K$ ,  $s_{\overline{w}s} = \partial_{\overline{w}} \partial_s K$ . (22)

Generalized Kähler property guarantees that the vanishing of the Christoffel symbols of  $CP_2$  (rather than those of 3-dimensional projection), which are of type  $\begin{cases} k \\ \varepsilon \ \overline{\varepsilon} \end{cases}$ .

$$\left\{ {k \atop \xi \ \overline{\xi}} \right\} = 0 \quad . \tag{23}$$

Here the coordinates of  $CP_2$  have been chosen in such a manner that three of them correspond to the coordinates of the projection and fourth coordinate is constant at the projection. The upper index k refers also to the  $CP_2$  coordinate, which is constant for the  $CP_2$  projection. If energy momentum tensor has only components of type  $T^{+-}$  and  $T^{w\overline{w}}$ , field equations are satisfied even when if non-diagonal Christoffel symbols of  $CP_2$  are present. The challenge is to discover solution ansatz, which guarantees this property of the energy momentum tensor.

A stronger variant of Kähler property would be that also  $s_{ss}$  vanishes so that the coordinate lines defined by s would define light like curves in  $CP_2$ . The topologization of the Kähler current however implies that  $CP_2$  projection is a projection of a 3-surface with strong Kähler property. Using  $(s, \xi, \overline{\xi}, S^-)$  as coordinates for the space-time surface defined by the ansatz ( $w = w(\xi, s), S^+ =$  $S^+(s)$ ) one finds that  $g_{ss}$  must be vanishing so that stronger variant of the Kähler property holds true for  $S^- = constant$  3-surfaces.

The topologization condition for the Kähler current can be solved completely generally in terms of the induced metric using  $(\xi, \overline{\xi}, s)$  and some coordinate of  $M_+^4$ , call it  $x^4$ , as space-time coordinates. Topologization boils down to the conditions

$$\partial_{\beta} (J^{\alpha\beta} \sqrt{g}) = 0 \text{ for } \alpha \in \{\xi, \overline{\xi}, s\} ,$$
  
$$g^{4i} \neq 0 .$$
(24)

Thus 3-dimensional empty space Maxwell equations and the non-orthogonality of  $X^4$  coordinate lines and the 3-surfaces defined by the lift of the  $CP_2$  projection.

#### **3.5.5** A solution ansatz yielding light-like current in $D_{CP_2} = 3$ case

The basic idea is that of generalized Kähler structure and solutions of field equations as maps or deformations of canonically imbedded  $M_+^4$  respecting this structure and guaranteing that the only non-vanishing components of the energy momentum tensor are  $T^{\xi\xi}$  and  $T^{s-}$  in the coordinates  $(\xi, \overline{\xi}, s, S^-)$ .

1. The coordinates  $(w, S^+)$  are assumed to holomorphic functions of the  $CP_2$  coordinates  $(s, \xi)$ 

$$S^+ = S^+(s) , \quad w = w(\xi, s) .$$
 (25)

Obviously  $S^+$  could be replaced with  $S^-$ . The ansatz is completely symmetric with respect to the exchange of the roles of (s, w) and  $(S^+, \xi)$  since it maps longitudinal degrees of freedom to longitudinal ones and transverse degrees of freedom to transverse ones.

- 2. Field equations are satisfied if the only non-vanishing components of the energy momentum tensor are of type  $T^{\xi\bar{\xi}}$  and  $T^{s-}$ . The reason is that the  $CP_2$  Christoffel symbols for projection and projections of  $M^4_+$  Christoffel symbols are vanishing for these lower index pairs.
- 3. By a straightforward calculation one can verify that the only manner to achieve the required structure of energy momentum tensor is to assume that the induced metric in the coordinates  $(\xi, \overline{\xi}, s, S^-)$  has as non-vanishing components only  $g_{\xi\overline{\xi}}$  and  $g_{s-}$

$$g_{ss} = 0$$
 ,  $g_{\xi s} = 0$  ,  $g_{\overline{\xi}s} = 0$  . (26)

Obviously the space-time surface must factorize into an orthogonal product of longitudinal and transversal spaces.

4. The condition guaranteing the product structure of the metric is

$$s_{ss} = m_{+w}\partial_s w(\xi, s)\partial_s S^+(s) + m_{+\overline{w}}\partial_s \overline{w}(\xi, s)\partial_s S^+(s) ,$$
  

$$s_{s\xi} = m_{+w}\partial_{\xi} w(\xi)\partial_s S^+(s) ,$$
  

$$s_{s\overline{\xi}} = m_{+w}\partial_{\overline{\xi}} w(\overline{\xi})\partial_s S^+(s) .$$
(27)

Thus the function of dynamics is to diagonalize the metric and provide it with strong Kähler property. Obviously the  $CP_2$  projection corresponds to a light-like surface for all values of  $S^-$  so that space-time surface is foliated by light-like surfaces and the notion of generalized conformal invariance makes sense for the entire space-time surface rather than only for its boundary or elementary particle horizons.

5. The requirement that the Kähler current is proportional to the instanton current means that only the  $j^-$  component of the current is non-vanishing. This gives the following conditions

$$j^{\xi}\sqrt{g} = \partial_{\beta}(J^{\xi\beta}\sqrt{g}) = 0 , \quad j^{\overline{\xi}}\sqrt{g} = \partial_{\beta}(J^{\overline{\xi}\beta}\sqrt{g}) = 0 ,$$
  
$$j^{+}\sqrt{g} = \partial_{\beta}(J^{+\beta}\sqrt{g}) = 0 .$$
 (28)

Since  $J^{+\beta}$  vanishes, the condition

$$\sqrt{g}j^+ = \partial_\beta (J^{+\beta}\sqrt{g}) = 0 \tag{29}$$

is identically satisfied. Therefore the number of field equations reduces to three.

The physical interpretation of the solution ansatz deserves some comments.

- 1. The light-like character of the Kähler current brings in mind  $CP_2$  extremals for which  $CP_2$  projection is light like. This suggests that the topological condensation of  $CP_2$  type extremal occurs on  $D_{CP_2} = 3$  helical space-time sheet representing zitterbewegung. In the case of many-body system light-likeness of the current does not require that particles are massless if particles of opposite charges can be present. Field tensor has the form  $(J^{\xi\bar{\xi}}, J^{\xi-}, J^{\bar{\xi}-})$ . Both helical magnetic field and electric field present as is clear when one replaces the coordinates  $(S^+, S^-)$  with time-like and space-like coordinate. Magnetic field dominates but the presence of electric field means that genuine Beltrami field is not in question.
- 2. Since the induced metric is product metric, 3-surface is metrically product of 2-dimensional surface  $X^2$  and line or circle and obeys product topology. If absolute minima correspond to asymptotic self-organization patterns, the appearance of the product topology and even metric is not so surprising. Thus the solutions can be classified by the genus of  $X^2$ . An interesting question is how closely the explanation of family replication phenomenon in terms of the topology of the boundary component of elementary particle like 3-surface relates to this. The heaviness and instability of particles which correspond to genera g > 2 (sphere with more than two handles) might have simple explanation as absence of (stable)  $D_{CP_2} = 3$  solutions of field equations with genus g > 2.

- 3. The solution ansatz need not be the most general. Kähler current is light-like and already this is enough to reduce the field equations to the form involving only energy momentum tensor. One might hope of finding also solution ansätze for which Kähler current is timelike or space-like. Space-likeness of the Kähler current might be achieved if the complex coordinates  $(\xi, \overline{\xi})$  and hyper-complex coordinates  $(S^+, S^-)$  change the role. For this solution ansatz electric field would dominate. Note that the possibility that Kähler current is always light-like cannot be excluded.
- 4. Suppose that  $CP_2$  projection quite generally defines a foliation of the space-time surface by light-like 3-surfaces, as is suggested by the conformal invariance. If the induced metric has Minkowskian signature, the fourth coordinate  $x^4$  and thus also Kähler current must be time-like or light-like so that magnetic field dominates. Already the requirement that the metric is non-degenerate implies  $g_{s4} \neq 0$  so that the metric for the  $\xi = constant$  2-surfaces has a Minkowskian signature. Thus space-like Kähler current does not allow the lift of the  $CP_2$  projection to be light-like.

# **3.5.6** Are solutions with time-like or space-like Kähler current possible in $D_{CP_2} = 3$ case?

The following ansatz gives good hopes for obtaining solutions with space-like and time-like Kähler currents.

- 1. Assign to light-like coordinates coordinates (T, Z) by the formula  $T = S^+ + S^-$  and  $Z = S^+ S^-$ . Space-time coordinates are taken to be  $(\xi, \overline{\xi}, s)$  and coordinate Z. The solution ansatz with time-like Kähler current results when the roles of T and Z are changed. It will however found that same solution ansatz can give rise to both space-like and time-like Kähler current.
- 2. The solution ansatz giving rise to a space-like Kähler current is defined by the equations

$$T = T(Z, s)$$
,  $w = w(\xi, s)$ . (30)

If T depends strongly on Z, the  $g_{ZZ}$  component of the induced metric becomes positive and Kähler current time-like.

3. The components of the induced metric are

$$g_{ZZ} = m_{ZZ} + m_{TT}\partial_{Z}T\partial_{s}T , \quad g_{Zs} = m_{TT}\partial_{Z}T\partial_{s}T ,$$

$$g_{ss} = s_{ss} + m_{TT}\partial_{s}T\partial_{s}T , \qquad g_{w\overline{w}} = s_{w\overline{w}} + m_{w\overline{w}}\partial_{\xi}w\partial_{\overline{\xi}}\overline{w} , \qquad (31)$$

$$g_{s\xi} = s_{s\xi} , \qquad \qquad g_{s\overline{\xi}} = s_{s\overline{\xi}} .$$

Topologized Kähler current has only Z-component and 3-dimensional empty space Maxwell's equations guarantee the topologization.

In  $CP_2$  degrees of freedom the contractions of the energy momentum tensor with Christoffel symbols vanish if  $T^{ss}$ ,  $T^{\xi s}$  and  $T^{\xi \xi}$  vanish as required by internal consistency. This is guaranteed if the condition

$$J^{\xi s} = 0 \tag{32}$$

holds true. Note however that  $J^{\xi Z}$  is non-vanishing. Therefore only the components  $T^{\xi \overline{\xi}}$  and  $T^{Z\xi}$ ,  $T^{Z\overline{\xi}}$  of energy momentum tensor are non-vanishing, and field equations reduce to the conditions

$$\partial_{\overline{\xi}}(J^{\xi\overline{\xi}}\sqrt{g}) + \partial_Z(J^{\xi Z}\sqrt{g}) = 0 ,$$
  
$$\partial_{\xi}(J^{\overline{\xi}\xi}\sqrt{g}) + \partial_Z(J^{\overline{\xi}Z}\sqrt{g}) = 0 .$$
(33)

In the special case that the induce metric does not depend on z-coordinate equations reduce to holomorphicity conditions. This is achieve if T depends linearly on Z: T = aZ.

The contractions with  $M_+^4$  Christoffel symbols come from the non-vanishing of  $T^{Z\xi}$  and vanish if the Hamilton Jacobi structure satisfies the conditions

$$\begin{cases} {}^{k}_{T w} \\ {}^{k}_{T w} \end{cases} = 0 , \quad \begin{cases} {}^{k}_{T \overline{w}} \\ {}^{k}_{Z w} \end{cases} = 0 , \quad \begin{cases} {}^{k}_{Z \overline{w}} \\ {}^{k}_{Z \overline{w}} \end{cases} = 0$$

$$(34)$$

hold true. The conditions are equivalent with the conditions

$${k \\ \pm w} = 0 , \quad {k \\ \pm \overline{w}} = 0 .$$
 (35)

These conditions possess solutions (standard light cone coordinates are the simplest example). Also the second derivatives of T(s, Z) contribute to the second fundamental form but they do not give rise to non-vanishing contractions with the energy momentum tensor. The cautious conclusion is that also solutions with time-like or space-like Kähler current are possible.

#### **3.5.7** $D_{CP_2} = 4$ case

The preceding discussion was for  $D_{CP_2} = 3$  and one should generalize the discussion to  $D_{CP_2} = 4$  case.

- 1. Hamilton Jacobi structure for  $M^4_+$  is expected to be crucial also now.
- 2. One might hope that for D = 4 the Kähler structure of  $CP_2$  defines a foliation of  $CP_2$  by 3-dimensional contact structures. This requires that there is a coordinate varying along the field lines of the normal vector field X defined as the dual of the three-form  $A \wedge dA = A \wedge J$ . By the previous considerations the condition for this reads as  $dX = d(\log \phi) \wedge X$  and implies  $X \wedge dX = 0$ . Using the self duality of the Kähler form one can express X as  $X^k = J^{kl}A_l$ . By a brief calculation one finds that  $X \wedge dX \propto X$  holds true so that (somewhat disappointingly) a foliation of  $CP_2$  by contact structures does not exist.

For  $D_{CP_2} = 4$  case Kähler current vanishes and this case corresponds to what I have called earlier Maxwellian phase since empty space Maxwell's equations are indeed satisfied.

1. Solution ansatz with a 3-dimensional  $M^4_+$  projection

The basic idea is that the complex structure of  $CP_2$  is preserved so that one can use complex coordinates  $(\xi^1, \xi^2)$  for  $CP_2$  in which  $CP_2$  Christoffel symbols and energy momentum tensor have automatically the desired properties. This is achieved the second light like coordinate, say v, is

non-dynamical so that the induced metric does not receive any contribution from the longitudinal degrees of freedom. In this case one has

$$S^+ = S^+(\xi^1, \xi^2)$$
,  $w = w(\xi^1, \xi^2)$ ,  $S^- = constant$ . (36)

The induced metric does possesses only components of type  $g_{i\bar{i}}$  if the conditions

$$g_{+w} = 0 , \quad g_{+\overline{w}} = 0 .$$
 (37)

This guarantees that energy momentum tensor has only components of type  $T^{i\bar{j}}$  in coordinates  $(\xi^1, \xi^2)$  and their contractions with the Christoffel symbols of  $CP_2$  vanish identically. In  $M_+^4$  degrees of freedom one must pose the conditions

$${k \atop w_+} = 0$$
,  ${k \atop \overline{w}_+} = 0$ ,  ${k \atop ++} = 0$ . (38)

on Christoffel symbols. These conditions are satisfied if the  $M^4_+$  metric does not depend on  $S^+$ :

$$\partial_+ m_{kl} = 0 . (39)$$

This means that  $m_{-w}$  and  $m_{-\overline{w}}$  can be non-vanishing but like  $m_{+-}$  they cannot depend on  $S^+$ .

The second derivatives of  $S^+$  appearing in the second fundamental form are also a source of trouble unless they vanish. Hence  $S^+$  must be a linear function of the coordinates  $\xi^k$ :

$$S^+ = a_k \xi^k + \overline{a}_k \overline{\xi}^k \quad . \tag{40}$$

Field equations are the counterparts of empty space Maxwell equations  $j^{\alpha} = 0$  but with  $M_{+}^{4}$  coordinates (u, w) appearing as dynamical variables and entering only through the induced metric. By holomorphy the field equations can be written as

$$\partial_j (J^{j\bar{i}} \sqrt{g}) = 0 , \quad \partial_{\bar{j}} (J^{\bar{j}i} \sqrt{g}) = 0 , \qquad (41)$$

and can be interpreted as conditions stating the holomorphy of the contravariant Kähler form.

What is remarkable is that the  $M^4_+$  projection of the solution is 3-dimensional light like surface and that the induced metric has Euclidian signature. Light front would become a concrete geometric object with one compactified dimension rather than being a mere conceptualization. One could see this as topological quantization for the notion of light front or of electromagnetic shock wave, or perhaps even as the realization of the particle aspect of gauge fields at classical level.

If the latter interpretation is correct, quantum classical correspondence would be realized very concretely. Wave and particle aspects would both be present. One could understand the interactions of charged particles with electromagnetic fields both in terms of absorption and emission of topological field quanta and in terms of the interaction with a classical field as particle topologically condenses at the photonic light front.

For  $CP_2$  type extremals for which  $M_+^4$  projection is a light like curve correspond to a special case of this solution ansatz: transversal  $M_+^4$  coordinates are constant and  $S^+$  is now arbitrary function of  $CP_2$  coordinates. This is possible since  $M_+^4$  projection is 1-dimensional.

2. Are solutions with a 4-dimensional  $M^4_+$  projection possible?

The most natural solution ansatz is the one for which  $CP_2$  complex structure is preserved so that energy momentum tensor has desired properties. For four-dimensional  $M_+^4$  projection this ansatz does not seem to make promising since the contribution of the longitudinal degrees of freedom implies that the induced metric is not anymore of desired form since the components  $g_{ij} = m_{+-}(\partial_{\xi^i}S^+\partial_{\xi^j}S^- + m_{+-}\partial_{\xi^i}S^-\partial_{\xi^j}S^+)$  are non-vanishing.

- 1. The natural dynamical variables are still Minkowski coordinates  $(w, \overline{w}, S^+, S^-)$  for some Hamilton Jacobi structure. Since the complex structure of  $CP_2$  must be given up,  $CP_2$ coordinates can be written as  $(\xi, s, r)$  to stress the fact that only "one half" of the Kähler structure of  $CP_2$  is respected by the solution ansatz.
- 2. The solution ansatz has the same general form as in D = 3 case and must be symmetric with respect to the exchange of  $M_{+}^{4}$  and  $CP_{2}$  coordinates. Transverse coordinates are mapped to transverse ones and longitudinal coordinates to longitudinal ones:

$$(S^+, S^-) = (S^+(s, r), S^-(s, r)) \quad , \quad w = w(\xi) \quad .$$
(42)

This ansatz would describe ordinary Maxwell field in  $M_+^4$  since the roles of  $M_+^4$  coordinates and  $CP_2$  coordinates are interchangeable.

It is however far from obvious whether there are any solutions with a 4-dimensional  $M_{+}^{4}$  projection. That empty space Maxwell's equations would allow only the topologically quantized light fronts as its solutions would realize quantum classical correspondence very concretely.

#### **3.5.8** $D_{CP_2} = 2$ case

Hamilton Jacobi structure for  $M_+^4$  is assumed also for  $D_{CP_2} = 2$ , whereas the contact structure for  $CP_2$  is in D = 2 case replaced by the induced Kähler structure. Topologization yields vanishing Kähler current. Light-likeness provides a second manner to achieve vanishing Lorentz force but one cannot exclude the possibility of time- and space-like Kähler current.

- 1. Solutions with vanishing Kähler current
- 1. String like objects, which are products X<sup>2</sup> × Y<sup>2</sup> ⊂ M<sup>4</sup><sub>+</sub> × CP<sub>2</sub> of minimal surfaces Y<sup>2</sup> of M<sup>4</sup><sub>+</sub> with geodesic spheres S<sup>2</sup> of CP<sub>2</sub> and carry vanishing gauge current. String like objects allow considerable generalization from simple Cartesian products of X<sup>2</sup> × Y<sup>2</sup> ⊂ M<sup>4</sup> × S<sup>2</sup>. Let (w, w, S<sup>+</sup>, S<sup>-</sup>) define the Hamilton Jacobi structure for M<sup>4</sup><sub>+</sub>. w = constant surfaces define minimal surfaces X<sup>2</sup> of M<sup>4</sup><sub>+</sub>. Let ξ denote complex coordinate for a sub-manifold of CP<sub>2</sub> such that the imbedding to CP<sub>2</sub> is holomorphic: (ξ<sup>1</sup>, ξ<sup>2</sup>) = (f<sup>1</sup>(ξ), f<sup>2</sup>(ξ)). The resulting surface Y<sup>2</sup> ⊂ CP<sub>2</sub> is a minimal surface and field equations reduce to the requirement that the Kähler current vanishes: ∂<sub>ξ</sub>(J<sup>ξξ</sup>√g<sub>2</sub>) = 0. One-dimensional strings are deformed to 3-dimensional cylinders representing magnetic flux tubes. The oscillations of string correspond to waves moving along string with light velocity, and for more general solutions they become TGD counterparts of Alfwen waves associated with magnetic flux tubes regarded as oscillations of magnetic flux lines behaving effectively like strings. It must be emphasized that Alfwen waves are a phenomenological notion not really justified by the properties of Maxwell's equations.
- 2. Also electret type solutions with the role of the magnetic field taken by the electric field are possible.  $(\xi, \overline{\xi}, u, v)$  would provide the natural coordinates and the solution ansatz would be of the form

$$(s,r) = (s(u,v), r(u,v)) \quad , \quad \xi = constant \quad , \tag{43}$$

and corresponds to a vanishing Kähler current.

3. Both magnetic and electric fields are necessarily present only for the solutions carrying nonvanishing electric charge density (proportional to  $\overline{B} \cdot \overline{A}$ ). Thus one can ask whether more general solutions carrying both magnetic and electric field are possible. As a matter fact, one must first answer the question what one really means with the magnetic field. By choosing the coordinates of 2-dimensional  $CP_2$  projection as space-time coordinates one can define what one means with magnetic and electric field in a coordinate invariant manner. Since the  $CP_2$  Kähler form for the  $CP_2$  projection with  $D_{CP_2} = 2$  can be regarded as a pure Kähler magnetic field, the induced Kähler field is either magnetic field or electric field.

The form of the ansatz would be

$$(s,r) = (s,r)(u,v,w,\overline{w}) , \quad \xi = constant .$$
(44)

As a matter fact,  $CP_2$  coordinates depend on two properly chosen  $M^4$  coordinates only.

1. Solutions with light-like Kähler current

There are large classes of solutions of field equations with a light-like Kähler current and 2dimensional  $CP_2$  projection.

- 1. Massless extremals for which  $CP_2$  coordinates are arbitrary functions of one transversal coordinate  $e = f(w, \overline{w})$  defining local polarization direction and light like coordinate u of  $M_+^4$  and carrying in the general case a light like current. In this case the holomorphy does not play any role.
- 2. The string like solutions thickened to magnetic flux tubes carrying TGD counterparts of Alfwen waves generalize to solutions allowing also light-like Kähler current. Also now Kähler metric is allowed to develop a component between longitudinal and transversal degrees of freedom so that Kähler current develops a light-like component. The ansatz is of the form

$$\xi^{i} = f^{i}(\xi)$$
,  $w = w(\xi)$ ,  $S^{-} = s^{-}$ ,  $S^{+} = s^{+} + f(\xi, \overline{\xi})$ .

Only the components  $g_{+\xi}$  and  $g_{+\overline{\xi}}$  of the induced metric receive contributions from the modification of the solution ansatz. The contravariant metric receives contributions to  $g^{-\xi}$  and  $g^{-\overline{\xi}}$  whereas  $g^{+\xi}$  and  $g^{+\overline{\xi}}$  remain zero. Since the partial derivatives  $\partial_{\xi}\partial_{+}h^{k}$  and  $\partial_{\overline{\xi}}\partial_{+}h^{k}$  and corresponding projections of Christoffel symbols vanish, field equations are satisfied. Kähler current develops a non-vanishing component  $j^{-}$ . Apart from the presence of the electric field, these solutions are highly analogous to Beltrami fields.

3. Do scalar wave pulses represent a solution type with non-vanishing but not light-like Kähler current?

Since longitudinal polarizations are possible only for off mass shell virtual photons, physical intuition suggests that scalar wave pulse solutions describing the propagation of longitudinal electric field with light velocity cannot appear as asymptotic field patterns. This is also consistent with the claim that scalar wave pulses are associated with the transients involved with sudden

switching of electric voltage on or off. Let  $M^4 = M^2 \oplus E^2$  be the standard decomposition of  $M^4$  to flat longitudinal and transversal spaces, and  $S^2$  a homologically non-trivial geodesic sphere of  $CP_2$ . The simplest solution ansatz corresponds to a surface  $X^2 \times Y^2$ ,  $X^2 \subset E^2$ , such that  $Y^2$  is a surface defined by a map  $S^2 \to M^2$  (or vice versa).

Energy momentum tensor is in both longitudinal and transversal degrees of freedom proportional to the corresponding part of the induced metric. Field equations are trivially true in the transversal degrees of freedom. The calculation of the divergence of energy momentum tensor demonstrates that Kähler current can be regarded as a vector field

$$j^{\alpha} = \frac{1}{4} J^{\alpha\beta} \partial_{\beta} L$$

defined by the Kähler action density acting as Hamiltonian. Poisson bracket is defined by the pseudo-symplectic form associated with the induced Kähler form with respect to the induced metric rather that that of  $S^2$  (using  $S^2$ -coordinates as coordinates for  $Y^2$ , the square of this pseudo-symplectic form is equal to metric multiplied by the ratio  $det(g(Y^2))/det(g(S^2)))$ .

In longitudinal degrees of freedom field equations are minimal surface equations with a source term proportional to the Kähler current divided by the Kähler action density. The vanishing of the Kähler current is possible only if Kähler action density is constant. This condition is true in the approximation that the induced metric for  $Y^2$  is flat, that is at the limit when  $M^4$  projection has size larger than size of  $CP_2$  projection and that induced metric has Minkowskian signature). It is not clear whether the minimal surface property of  $Y^2$  in  $M^2 \times S^2$  is consistent with the constancy of the Kähler action density. This would suggest that classical gravitational interactions eliminate scalar wave pulses as asymptotic field patterns and cause the deviation from the minimal surface property and the non-vanishing of the Kähler current. The fact that solution becomes "instanton" like Euclidian solution when  $S^+$  and  $S^-$  become constant suggests that the  $M^4$  projection of the solution quite generally has a finite extension in time direction.

### **3.5.9** Could $D_{CP_2} = 2 \rightarrow 3$ transition occur in rotating magnetic systems?

I have studied the imbeddings of simple cylindrical and helical magnetic fields in various applications of TGD to condensed matter systems, in particular in attempts to understand the strange findings about rotating magnetic systems [G2].

Let  $S^2$  be the homologically non-trivial geodesic sphere of  $CP_2$  with standard spherical coordinates  $(U \equiv cos(\theta), \Phi)$  and let  $(t, \rho, \phi, z)$  denote cylindrical coordinates for a cylindrical space-time sheet. The simplest possible space-time surfaces  $X^4 \subset M^4_+ \times S^2$  carrying helical Kähler magnetic field depending on the radial cylindrical coordinate  $\rho$ , are given by:

$$U = U(\rho) , \qquad \Phi = n\phi + kz ,$$
  

$$J_{\rho\phi} = n\partial_{\rho}U , \qquad J_{\rho z} = k\partial_{\rho}U .$$
(45)

This helical field is not Beltrami field as one can easily find. A more general ansatz corresponding defined by

$$\Phi = \omega t + kz + n\phi$$

would in cylindrical coordinates give rise to both helical magnetic field and radial electric field depending on  $\rho$  only. This field can be obtained by simply replacing the vector potential with its rotated version and provides the natural first approximation for the fields associated with rotating magnetic systems.

A non-vanishing vacuum charge density is however generated when a constant magnetic field is put into rotation and is implied by the condition  $\overline{E} = \overline{v} \times \overline{B}$  stating vanishing of the Lorentz force. This condition does not follow from the induction law of Faraday although Faraday observed this effect first. This is also clear from the fact that the sign of the charge density depends on the direction of rotation.

The non-vanishing charge density is not consistent with the vanishing of the Kähler 4-current and requires a 3-dimensional  $CP_2$  projection and topologization of the Kähler current. Beltrami condition cannot hold true exactly for the rotating system. The conclusion is that rotation induces a phase transition  $D_{CP_2} = 2 \rightarrow 3$ . This could help to understand various strange effects related to the rotating magnetic systems [G2]. For instance, the increase of the dimension of  $CP_2$  projection could generate join along boundaries contacts and wormhole contacts leading to the transfer of charge between different space-time sheets. The possibly resulting flow of gravitational flux to larger space-time sheets might help to explain the claimed antigravity effects.

## **3.6** D = 3 phase allows infinite number of topological charges characterizing the linking of magnetic field lines

When space-time sheet possesses a D = 3-dimensional  $CP_2$  projection, one can assign to it a nonvanishing and conserved topological charge characterizing the linking of the magnetic field lines defined by Chern-Simons action density  $A \wedge dA/4\pi$  for induced Kähler form. This charge can be seen as classical topological invariant of the linked structure formed by magnetic field lines.

The topological charge can also vanish for D = 3 space-time sheets. In Darboux coordinates for which Kähler gauge potential reads as  $A = P_k dQ^k$ , the surfaces of this kind result if one has  $Q^2 = f(Q^1)$  implying  $A = f dQ^1$ ,  $f = P_1 + P_2 \partial_{Q_1} Q^2$ , which implies the condition  $A \wedge dA = 0$ . For these space-time sheets one can introduce  $Q^1$  as a global coordinate along field lines of A and define the phase factor  $exp(i \int A_\mu dx^\mu)$  as a wave function defined for the entire space-time sheet. This function could be interpreted as a phase of an order order parameter of super-conductor like state and there is a high temptation to assume that quantum coherence in this sense is lost for more general D = 3 solutions.

Chern-Simons action is known as helicity in electrodynamics [22]. Helicity indeed describes the linking of magnetic flux lines as is easy to see by interpreting magnetic field as incompressible fluid flow having A as vector potential:  $B = \nabla \times A$ . One can write A using the inverse of  $\nabla \times$  as  $A = (1/\nabla \times)B$ . The inverse is non-local operator expressible as

$$\frac{1}{\nabla \times} B(r) = \int dV' \frac{(r-r')}{|r-r'|^3} \times B(r') \quad ,$$

as a little calculation shows. This allows to write  $\int A \cdot B$  as

$$\int dV A \cdot B = \int dV dV' B(r) \cdot \left(\frac{(r-r')}{|r-r'|^3} \times B(r')\right) \quad ,$$

which is completely analogous to the Gauss formula for linking number when linked curves are replaced by a distribution of linked curves and an average is taken.

For D = 3 field equations imply that Kähler current is proportional to the helicity current by a factor which depends on  $CP_2$  coordinates, which implies that the current is automatically divergence free and defines a conserved charge for D = 3-dimensional  $CP_2$  projection for which the instanton density vanishes identically. Kähler charge is not equal to the helicity defined by the inner product of magnetic field and vector potential but to a more general topological charge.

The number of conserved topological charges is infinite since the product of any function of  $CP_2$ coordinates with the helicity current has vanishing divergence and defines a topological charge. A very natural function basis is provided by the scalar spherical harmonics of SU(3) defining Hamiltonians of  $CP_2$  canonical transformations and possessing well defined color quantum numbers. These functions define and infinite number of conserved charges which are also classical knot invariants in the sense that they are not affected at all when the 3-surface interpreted as a map from  $CP_2$ projection to  $M_+^4$  is deformed in  $M_+^4$  degrees of freedom. Also canonical transformations induced by Hamiltonians in irreducible representations of color group affect these invariants via Poisson bracket action when the U(1) gauge transformation induced by the canonical transformation corresponds to a single valued scalar function. These link invariants are additive in union whereas the quantum invariants defined by topological quantum field theories are multiplicative.

Also non-Abelian topological charges are well-defined. One can generalize the topological current associated with the Kähler form to a corresponding current associated with the induced electro-weak gauge fields whereas for classical color gauge fields the Chern-Simons form vanishes identically. Also in this case one can multiply the current by  $CP_2$  color harmonics to obtain an infinite number of invariants in D = 3 case. The only difference is that  $A \wedge dA$  is replaced by  $Tr(A \wedge (dA + 2A \wedge A/3))$ .

There is a strong temptation to assume that these conserved charges characterize colored quantum states of the conformally invariant quantum theory as a functional of the light-like 3-surface defining boundary of space-time sheet or elementary particle horizon surrounding wormhole contacts. They would be TGD analogs of the states of the topological quantum field theory defined by Chern-Simons action as highest weight states associated with corresponding Wess-Zumino-Witten theory. These charges could be interpreted as topological counterparts of the isometry charges of configuration space of 3-surfaces defined by the algebra of canonical transformations of  $CP_2$ .

The interpretation of these charges as contributions of light-like boundaries to configuration space Hamiltonians would be natural. The dynamics of the induced second quantized spinor fields relates to that of Kähler action by a super-symmetry, so that it should define super-symmetric counterparts of these knot invariants. The anti-commutators of these super charges cannot however contribute to configuration space Kähler metric so that topological zero modes are in question. These Hamiltonians and their super-charge counterparts would be responsible for the topological sector of quantum TGD.

# 3.7 Is the principle selecting preferred extremals of Kähler action equivalent with the topologization/light-likeness of Kähler current and with second law?

The basic question is whether the Kähler current is either topologized or light-like for all extremals or only for the preferred extremals of Kähler action analogous to Bohr orbits in some sense, presumably asymptotically as suggested by the fact that generalized Beltrami fields correspond to asymptotic self-organization patterns, when dissipation has become insignificant.

- 1. The generalized Beltrami conditions or light-likeness can hold true only asymptotically. First of all, generic non-asymptotic field configurations have  $D_{CP_2} = 4$ , and would thus carry a vanishing Kähler four-current if Beltrami conditions were satisfied universally rather than only asymptotically.  $j^{\alpha} = 0$  would obviously hold true also for the asymptotic configurations, in particular those with  $D_{CP_2} < 4$  so that empty space Maxwell's field equations would be universally satisfied for asymptotic field configurations with  $D_{CP_2} < 4$ .
- 2. The failure of the generalized Beltrami conditions would mean that Kähler field is completely analogous to a dissipative Maxwell field since  $\overline{j} \cdot \overline{E}$  is non-vanishing (note that isometry currents are conserved although energy momentum tensor is not). Quantum classical correspondence states that classical space-time dynamics is by its classical non-determinism able to mimic the non-deterministic sequence of quantum jumps at space-time level, in particular dissipation in various length scales defined by the hierarchy of space-time sheets. Classical fields could represent "symbolically" the average dynamics, in particular dissipation, in

shorter length scales. For instance, vacuum 4-current would be a symbolic representation for the average of the currents consisting of elementary particles.

The obvious objection to the idea is that second law realized as an asymptotic vanishing of Lorentz-Kähler force implies that all 3-surfaces approaching same asymptotic state have the same value of Kähler function. This is actually not a problem since it means an additional symmetry extending general coordinate invariance. The exponent of Kähler function would be highly analogous to a partition function defined as an exponent of Hamiltonian with Kähler coupling strength playing the role of temperature. The principle selecting preferred extremals of Kähler action would guarantee that the predicted both signs of Kähler coupling strength and depending on sign Kähler magnetic or electric field configurations are stable. The two phases would relate by time reversal to each other. The role of time reversed dissipation manifesting itself as processes like self assembly would involve in essential manner the time reversed negative energy world. The fact that Kähler coupling strength is of opposite sign for the time reversed dynamics is essential for internal consistency. For instance, creation of matter from vacuum in big bang can be seen as time reflection of stable negative energy cosmic strings as positive energy cosmic strings unstable against decay to magnetic flux tubes.

#### 3.7.1 Is the principle selecting preferred extremals equivalent with generalized Beltrami conditions?

Previous findings inspire the hypothesis that generalized Beltrami conditions express algebraically the conditions selecting preferred extremal of Kähler action so that they make sense also in the p-adic case.

- 1. Generalized Beltrami conditions are satisfied by the asymptotic field configurations representing self-organization patterns. For non-asymptotic fields vacuum Lorentz force is nonvanishing and does work in Maxwellian sense so that  $\overline{j} \cdot \overline{E}$  is non-vanishing. This would mean that the dynamics defined by Kähler action could in principle predict even the values of the parameters related to dissipation such as conductivities and viscosities. The space-time sheets of the many-sheeted space-time would be busily modelling its own physics in shorter length scales.
- 2. Absolute minimization of Kähler action implies that single space-time surface goes through given 3-surface apart from the non-uniqueness caused by the non-determinism of Kähler action. This gives *four* additional local conditions to the initial values of field equations fixing the time derivatives of the four dynamical imbedding space coordinates (conditions are analogous to Bohr conditions).

The topologization of the Kähler current current gives also *four* local conditions:

i) For  $D_{CP_2} < 4$  the vanishing of instanton density gives one condition, and the proportionality of the Kähler current to instanton current gives 3 conditions since the proportionality factor is an arbitrary function of  $CP_2$  coordinates. Altogether this makes four conditions. ii) For  $D_{CP_2=4}$  the vanishing of the Kähler current gives four conditions.

This encourages to think that the principle selecting preferred extremals of Kähler action forces the asymptotic behavior (final values instead of initial values) to correspond to dissipation-less state characterized by the generalized Beltrami conditions.

3. Absolute minimization in a strict sense of the word does not make sense in the p-adic context since p-adic numbers are not well-ordered, and one cannot even define the action integral as a p-adic number except perhaps by algebraic continuation procedure described in the first part of the book. The generalized Beltrami conditions are however purely algebraic conditions and make sense also in the p-adic context. Therefore it possible to give a precise content to the notion of the principle selecting preferred extremals also in p-adic context.

#### 3.7.2 Is the principle selecting preferred extremals equivalent with the second law?

The fact that Beltrami conditions are associated with the asymptotic dynamics suggests that absolute minimization (or its variant discussed in [E2]) is equivalent with the second law at space-time level. Or putting it more cautiously: second law at space-time level could be equivalent with the principle selecting preferred extremals of Kähler action. If not, one must give up the principle selecting preferred extremals and replace it with the second law.

For space-time sheets with negative time orientation and negative energy, say "massless extremals" representing phase conjugate laser waves, field configurations would approach non-dissipating ones in the geometric past, and the arrow of geometric time would be opposite to the standard one in this case. This situation is possible for space-time sheets of finite duration, in particular virtual particle like space-time sheets or the negative energy space-time sheets extending down to the boundary of imbedding space (moment of "big bang"). This would explain at the space-time level the change of arrow of time and breaking of the second law observed for the phase conjugate laser waves (used to generate healing and error correction for instance). In TGD framework second law is not a producer of a thermal chaos but Darwinian selector since state function reduction and state preparation by self measurements lead from a state with positive entanglement entropy to that with a negative entanglement entropy (defined number theoretically), and possessing only finitely extended rational entanglement identifiable as a bound state entanglement.

Absolute minimization of Kähler action indeed induces long range correlations since the positive Kähler action of space-time sheets carrying magnetic fields must be compensated by the negative Kähler action of space-time sheets dominated by Kähler electric fields. The resulting non-local long range correlations could serve as correlates for bound state entanglement. More concretely, the stable join along boundaries bonds would be the correlates for bound state entanglement whereas topological light rays analogous to the exchange of virtual photons could serve as classical correlates for unbound entanglement. The closedness (periodicity) of the field lines of Beltrami fields for space-like Kähler current and periodicity of the field pattern for the time like Kähler current could be space-time correlates for the rational entanglement. The pinary expansions of rational numbers which are periodic after finite number of pinary digits indeed represent closed orbits in the set of integers modulo p. Amusingly, the first non-periodic pits of the expansion would in fact be analogous to the dissipative period.

Macro-temporal quantum coherence integrates sequences of quantum jumps to single effective quantum jump so that effectively a fractal hierarchy of quantum jumps emerges having the fractal hierarchy of time scales of dissipation resulting from many-sheetedness as a correlate. Even the anatomy of quantum jump could have space-time correlate. The final state of the quantum jump would correspond to highly negentropic and non-dissipating topologically quantized generalized Beltrami fields. State function reduction and preparation would correspond to the non-deterministic dissipative approach to the non-dissipative Beltrami field configuration. The points of space-time sheets with vanishing Kähler 4-currents would be unstable against quantum jumps generating an instability of the Beltrami field leading to a field configuration with a nonvanishing Lorentz 4-force and emission of topological light rays representing unstable entanglement. Quantum jump would have this kind of instability as a natural space-time correlate.

To sum up, the main lessons would be following.

- 1. The ability of basically non-dissipative dynamics to mimic dissipative dynamics in terms of energy momentum tensor would be the basic reason for why space-times must be 4-surfaces.
- 2. If absolute minimization or its variant discussed in [E2] of the Kähler action is correct principle, it must correspond to the second law, which is the Darwinian selector of the most

information rich patterns rather than a thermal killer.

#### 3.8 Generalized Beltrami fields and biological systems

The following arguments support the view that generalized Beltrami fields play a key role in living systems, and that  $D_{CP_2} = 2$  corresponds to ordered phase,  $D_{CP_2} = 3$  to spin glass phase and  $D_{CP_2} = 4$  to chaos, with  $D_{CP_2} = 3$  defining life as a phenomenon at the boundary between order and chaos.

#### 3.8.1 Why generalized Beltrami fields are important for living systems?

Chirality, complexity, and high level of organization make  $D_{CP_2} = 3$  generalized Beltrami fields excellent candidates for the magnetic bodies of living systems.

- 1. Chirality selection is one of the basic signatures of living systems. Beltrami field is characterized by a chirality defined by the relative sign of the current and magnetic field, which means parity breaking. Chirality reduces to the sign of the function  $\psi$  appearing in the topologization condition and makes sense also for the generalized Beltrami fields.
- 2. Although Beltrami fields can be extremely complex, they are also extremely organized. The reason is that the function  $\alpha$  is constant along flux lines so that flux lines must in the case of compact Riemann 3-manifold belong to 2-dimensional  $\alpha = constant$  closed surfaces, in fact two-dimensional invariant tori [20].

For generalized Beltrami fields the function  $\psi$  is constant along the flow lines of the Kähler current. Space-time sheets with 3-dimensional  $CP_2$  projection serve as an illustrative example. One can use the coordinates for the  $CP_2$  projection as space-time coordinates so that one space-time coordinate disappears totally from consideration. Hence the situation reduces to a flow in a 3-dimensional submanifold of  $CP_2$ . One can distinguish between three types of flow lines corresponding to space-like, light-like and time-like topological current. The 2-dimensional  $\psi = constant$  invariant manifolds are sub-manifolds of  $CP_2$ . Ordinary Beltrami fields are a special case of space-like flow with flow lines belonging to the 2-dimensional invariant tori of  $CP_2$ . Time-like and light-like situations are more complex since the flow lines need not be closed so that the 2-dimensional  $\psi = constant$ surfaces can have boundaries.

For periodic self-organization patterns flow lines are closed and  $\psi = constant$  surfaces of  $CP_2$ must be invariant tori. The dynamics of the periodic flow is obtained from that of a steady flow by replacing one spatial coordinate with effectively periodic time coordinate. Therefore topological notions like helix structure, linking, and knotting have a dynamical meaning at the level of  $CP_2$ projection. The periodic generalized Beltrami fields are highly organized also in the temporal domain despite the potentiality for extreme topological complexity.

For these reasons topologically quantized generalized Beltrami fields provide an excellent candidate for a generic model for the dynamics of biological self-organization patterns. A natural guess is that many-sheeted magnetic and  $Z^0$  magnetic fields and their generalizations serve as templates for the helical molecules populating living matter, and explain both chiral selection, the complex linking and knotting of DNA and protein molecules, and even the extremely complex and self-organized dynamics of biological systems at the molecular level.

The intricate topological structures of DNA, RNA, and protein molecules are known to have a deep significance besides their chemical structure, and they could even define something analogous to the genetic code. Usually the topology and geometry of bio-molecules is believed to reduce to chemistry. TGD suggests that space-like generalized Beltrami fields serve as templates for the formation of bio-molecules and bio-structures in general. The dynamics of bio-systems would in turn utilize the time-like Beltrami fields as templates. There could even exist a mapping from the

topology of magnetic flux tube structures serving as templates for bio-molecules to the templates of self-organized dynamics. The helical structures, knotting, and linking of bio-molecules would thus define a symbolic representation, and even coding for the dynamics of the bio-system analogous to written language.

#### **3.8.2** $D_{CP_2} = 3$ systems as boundary between $D_{CP_2} = 2$ order and $D_{CP_2} = 4$ chaos

The dimension of  $CP_2$  projection is basic classifier for the asymptotic self-organization patterns.

#### 1. $D_{CP_2} = 4$ phase, dead matter, and chaos

 $D_{CP_2} = 4$  corresponds to the ordinary Maxwellian phase in which Kähler current and charge density vanish and there is no topologization of Kähler current. By its maximal dimension this phase would naturally correspond to disordered phase, ordinary dead matter. If one assumes that Kähler charge corresponds to either em charge or  $Z^0$  charge then the signature of this state of matter would be em neutrality or  $Z^0$  neutrality.

2.  $D_{CP_2} = 2$  phase as ordered phase

By the low dimension of  $CP_2$  projection  $D_{CP_2} = 2$  phase is the least stable phase possible only at cold space-time sheets. Kähler current is either vanishing or light-like, and Beltrami fields are not possible. This phase is highly ordered and much like a topological quantized version of ferro-magnet. In particular, it is possible to have a global coordinate varying along the field lines of the vector potential also now. The magnetic and  $Z^0$  magnetic body of any system is a candidate for this kind of system.  $Z^0$  field is indeed always present for vacuum extremals having D = 2 and the vanishing of em field requires that that  $sin^2(\theta_W)$  ( $\theta_W$  is Weinberg angle) vanishes.

#### 3. $D_{CP_2} = 3$ corresponds to living matter

 $D_{CP_2} = 3$  corresponds to highly organized phase characterized in the case of space-like Kähler current by complex helical structures necessarily accompanied by topologized Kähler charge density  $\propto \overline{A} \cdot \overline{B} \neq 0$  and Kähler current  $\overline{E} \times \overline{A} + \phi \overline{B}$ . For time like Kähler currents the helical structures are replaced by periodic oscillation patterns for the state of the system. By the non-maximal dimension of  $CP_2$  projection this phase must be unstable against too strong external perturbations and cannot survive at too high temperatures. Living matter is thus excellent candidate for this phase and it might be that the interaction of the magnetic body with living matter makes possible the transition from  $D_{CP_2} = 2$  phase to the self-organizing  $D_{CP_2} = 3$  phase.

Living matter which is indeed populated by helical structures providing examples of space-like Kähler current. Strongly charged lipid layers of cell membrane might provide example of time-like Kähler current. Cell membrane, micro-tubuli, DNA, and proteins are known to be electrically charged and  $Z^0$  charge plays key role in TGD based model of catalysis discussed in [N4]. For instance, denaturing of DNA destroying its helical structure could be interpreted as a transition leading from D = 3 phase to D = 4 phase. The prediction is that the denatured phase should be electromagnetically (or  $Z^0$ ) neutral.

Beltrami fields result when Kähler charge density vanishes. For these configurations magnetic field and current density take the role of the vector potential and magnetic field as far as the contact structure is considered. For Beltrami fields there exist a global coordinate along the field lines of the vector potential but not along those of the magnetic field. As a consequence, the covariant consistency condition  $(\partial_s - qeA_s)\Psi = 0$  frequently appearing in the physics of super conducting systems would make sense along the flow lines of the vector potential for the order parameter of Bose-Einstein condensate. If Beltrami phase is super-conducting, then the state of the system must change in the transition to a more general phase. Since the field lines of the vector potential define chaotic orbits in this phase, the loss of coherence of the order parameter implying the loss of superconductivity by random collisions of particles is what one expects to happen. The existence of these three phases brings in mind systems allowing chaotic de-magnetized phase above critical temperature  $T_c$ , spin glass phase at the critical point, and ferromagnetic phase below  $T_c$ . Similar analogy is provided by liquid phase, liquid crystal phase possible in the vicinity of the critical point for liquid to solid transition, and solid phase. Perhaps one could regard  $D_{CP_2} = 3$  phase and life as a boundary region between  $D_{CP_2} = 2$  order and  $D_{CP_2} = 4$  chaos. This would naturally explain why life as it is known is possible in relatively narrow temperature interval.

# 4 Basic extremals of Kähler action

The solutions of field equations can be divided to vacuum extremals and non-vacuum extremal. Vacuum extremals come as two basic types:  $CP_2$  type vacuum extremals for which the induced Kähler field and Kähler action are non-vanishing and the extremals for which the induced Kähler field vanishes. The deformations of both extremals are expected to be of fundamental importance in TGD universe.

#### 4.1 *CP*<sub>2</sub> type vacuum extremals

These extremals correspond to various isometric imbeddings of  $CP_2$  to  $M_+^4 \times CP_2$ . One can also drill holes to  $CP_2$ . Using the coordinates of  $CP_2$  as coordinates for  $X^4$  the imbedding is given by the formula

$$m^{k} = m^{k}(u) ,$$
  

$$m_{kl}\dot{m}^{k}\dot{m}^{l} = 0 , \qquad (46)$$

where  $u(s^k)$  is an arbitrary function of  $CP_2$  coordinates. The latter condition tells that the curve representing the projection of  $X^4$  to  $M^4$  is light like curve. One can choose the functions  $m^i, i = 1, 2, 3$  freely and solve  $m^0$  from the condition expressing light likeness so that the number of this kind of extremals is very large.

The induced metric and Kähler field are just those of  $CP_2$  and energy momentum tensor  $T^{\alpha\beta}$ vanishes identically by the self duality of the Kähler form of  $CP_2$ . Also the canonical current  $j^{\alpha} = D_{\beta}J^{\alpha\beta}$  associated with the Kähler form vanishes identically. Therefore the field equations in the interior of  $X^4$  are satisfied. The field equations are also satisfied on the boundary components of  $CP_2$  type extremal because the non-vanishing boundary term is, besides the normal component of Kähler electric field, also proportional to the projection operator to the normal space and vanishes identically since the induced metric and Kähler form are identical with the metric and Kähler form of  $CP_2$ .

As a special case one obtains solutions for which  $M^4$  projection is light like geodesic. The projection of  $m^0 = constant$  surfaces to  $CP_2$  is u = constant 3-sub-manifold of  $CP_2$ . Geometrically these solutions correspond to a propagation of a massless particle. In a more general case the interpretation as an orbit of a massless particle is not the only possibility. For example, one can imagine a situation, where the center of mass of the particle is at rest and motion occurs along a circle at say  $(m^1, m^2)$  plane. The interpretation as a massive particle is natural. Amusingly, there is nice analogy with the classical theory of Dirac electron: massive Dirac fermion moves also with the velocity of light (zitterbewegung). The quantization of this random motion with light velocity leads to Virasoro conditions and this led to a breakthrough in the understanding of the symmetries of TGD. Super Virasoro invariance is a general symmetry of the configuration space geometry and quantum TGD.

The action for all extremals is same and given by the Kähler action for the imbedding of  $CP_2$ . The value of the action is given by

$$S = -\frac{\pi}{8\alpha_K} . \tag{47}$$

To derive this expression we have used the result that the value of Lagrangian is constant:  $L = 4/R^4$ , the volume of  $CP_2$  is  $V(CP_2) = \pi^2 R^4/2$  and the definition of the Kähler coupling strength  $k_1 = 1/16\pi\alpha_K$  (by definition,  $\pi R$  is the length of  $CP_2$  geodesics). Four-momentum vanishes for these extremals so that they can be regarded as vacuum extremals. The value of the action is negative so that these vacuum extremals are indeed favored by the minimization of the Kähler action. The the principle selecting preferred extremals of Kähler action suggests that ordinary vacuums with vanishing Kähler action density are unstable against the generation of  $CP_2$  type extremals. There are even reasons to expect that  $CP_2$  type extremals are for TGD what black holes are for GRT. Indeed, the nice generalization of the area law for the entropy of black hole [E5] supports this view.

In accordance with the basic ideas of TGD topologically condensed vacuum extremals should somehow correspond to massive particles. The properties of the  $CP_2$  type vacuum extremals are in accordance with this interpretation. Although these objects move with a velocity of light, the motion can be transformed to a mere so that the center of mass motion is trivial. Even the generation of the rest mass could might be understood classically as a consequence of the minimization of action. Long range Kähler fields generate negative action for the topologically condensed vacuum extremal (momentum zero massless particle) and Kähler field energy in turn is identifiable as the rest mass of the topologically condensed particle.

An interesting feature of these objects is that they can be regarded as gravitational instantons [23]. A further interesting feature of  $CP_2$  type extremals is that they carry nontrivial classical color charges. The possible relationship of this feature to color confinement raises interesting questions. Could one model classically the formation of the color singlets to take place through the emission of "colorons": states with zero momentum but non-vanishing color? Could these peculiar states reflect the infrared properties of the color interactions?

### 4.2 Vacuum extremals with vanishing Kähler field

Vacuum extremals correspond to 4-surfaces with vanishing Kähler field and therefore to gauge field zero configurations of gauge field theory. These surfaces have  $CP_2$  projection, which is Lagrange manifold. The condition expressing Lagrange manifold property is obtained in the following manner. Kähler potential of  $CP_2$  can be expressed in terms of the canonical coordinates  $(P_i, Q_i)$  for  $CP_2$  as

$$A = \sum_{k} P_k dQ^k \quad . \tag{48}$$

The conditions

$$P_k = \partial_{Q^k} f(Q^i) , \qquad (49)$$

where  $f(Q^i)$  is arbitrary function of its arguments, guarantee that Kähler potential is pure gauge. It is clear that canonical transformations, which act as local U(1) gauge transformations, transform different vacuum configurations to each other so that vacuum degeneracy is enormous. Also  $M_+^4$ diffeomorphisms act as the dynamical symmetries of the vacuum extremals. Some sub-group of these symmetries extends to the isometry group of the configuration space in the proposed construction of the configuration space metric. The vacuum degeneracy is still enhanced by the fact that the topology of the four-surface is practically free.

Vacuum extremals are certainly not absolute minima of the action. For the induced metric having Minkowski signature the generation of Kähler electric fields lowers the action. For Euclidian signature both electric and magnetic fields tend to reduce the action. Therefore the generation of Euclidian regions of space-time is expected to occur.  $CP_2$  type extremals, identifiable as real (as contrast to virtual) elementary particles, can be indeed regarded as these Euclidian regions.

Particle like vacuum extremals can be classified roughly by the number of the compactified dimensions D having size given by  $CP_2$  length. Thus one has D = 3 for  $CP_2$  type extremals, D = 2 for string like objects, D = 1 for membranes and D = 0 for pieces of  $M^4$ . As already mentioned, the rule  $h_{vac} = -D$  relating the vacuum weight of the Super Virasoro representation to the number of compactified dimensions of the vacuum extremal is very suggestive. D < 3 vacuum extremals would correspond in this picture to virtual particles, whose contribution to the generalized Feynmann diagram is not suppressed by the exponential of Kähler action unlike that associated with the virtual  $CP_2$  type lines.

 $M^4$  type vacuum extremals (representable as maps  $M^4_+ \to CP_2$  by definition) are also expected to be natural idealizations of the space-time at long length scales obtained by smoothing out small scale topological inhomogenities (particles) and therefore they should correspond to space-time of GRT in a reasonable approximation.

The reason would be "Yin-Yang principle" discussed in [D1].

- 1. Consider first the option for which Kähler function corresponds to an absolute minimum of Kähler action. Vacuum functional as an exponent of Kähler function is expected to concentrate on those 3-surfaces for which the Kähler action is non-negative. On the other hand, the requirement that Kähler action is absolute minimum for the space-time associated with a given 3-surface, tends to make the action negative. Therefore the vacuum functional is expected to differ considerably from zero only for 3-surfaces with a vanishing Kähler action per volume. It could also occur that the degeneracy of 3-surfaces with same large negative action compensates the exponent of Kähler function.
- 2. If preferred extrema correspond to Kähler calibrations or their duals [E2], Yin-Yang principle is modified to a more local principle. For Kähler calibrations (their duals) the absolute value of action in given region is minimized (maximized). A given region with a positive (negative sign) of action density favors Kähler electric (magnetic) fields. In long length scales the average density of Kähler action per four-volume tends to vanish so that Kähler function of the entire universe is expected to be very nearly zero. This regularizes the theory automatically and implies that average Kähler action per volume vanishes. Positive and finite values of Kähler function are of course favored.

In both cases the vanishing of Kähler action per volume in long length scales makes vacuum extremals excellent idealizations for the smoothed out space-time surface. Robertson-Walker cosmologies provide a good example in this respect. As a matter fact the smoothed out space-time is not a mere fictive concept since larger space-time sheets realize it as a essential part of the Universe.

Several absolute minima could be possible and the non-determinism of the vacuum extremals is not expected to be reduced completely. The remaining degeneracy could be even infinite. A good example is provided by the vacuum extremals representable as maps  $M_+^4 \rightarrow D^1$ , where  $D^1$  is one-dimensional curve of  $CP_2$ . This degeneracy could be interpreted as a space-time correlate for the non-determinism of quantum jumps with maximal deterministic regions representing quantum states in a sequence of quantum jumps.

#### 4.3 Cosmic strings

Cosmic strings are extremals of type  $X^2 \times S^2$ , where  $X^2$  is minimal surface in  $M^4_+$  (analogous to the orbit of a bosonic string) and  $S^2$  is the homologically non-trivial geodesic sphere of  $CP_2$ . The action of these extremals is positive and thus absolute minima are certainly not in question. One can however consider the possibility that these extremals are building blocks of the absolute minimum space-time surfaces since the principle selecting preferred extremals of the Kähler action is global rather than a local. Cosmic strings can contain also Kähler charged matter in the form of small holes containing elementary particle quantum numbers on their boundaries and the negative Kähler electric action for a topologically condensed cosmic string could cancel the Kähler magnetic action.

The string tension of the cosmic strings is given by

$$T = \frac{1}{8\alpha_K R^2} \simeq .2210^{-6} \frac{1}{G} \quad , \tag{50}$$

where  $\alpha_K \simeq \alpha_{em}$  has been used to get the numerical estimate. The string tension is of the same order of magnitude as the string tension of the cosmic strings of GUTs and this leads to the model of the galaxy formation providing a solution to the dark matter puzzle as well as to a model for large voids as caused by the presence of a strongly Kähler charged cosmic string. Cosmic strings play also fundamental role in the TGD inspired very early cosmology.

#### 4.4 Massless extremals

Massless extremals are characterized by massless wave vector p and polarization vector  $\varepsilon$  orthogonal to this wave vector. Using the coordinates of  $M^4$  as coordinates for  $X^4$  the solution is given as

$$s^{k} = f^{k}(u, v) ,$$
  

$$u = p \cdot m , \qquad v = \varepsilon \cdot m ,$$
  

$$p \cdot \varepsilon = 0 , \qquad p^{2} = 0 .$$
(51)

 $CP_2$  coordinates are arbitrary functions of  $p \cdot m$  and  $\varepsilon \cdot m$ . Clearly these solutions correspond to plane wave solutions of gauge field theories. It is important to notice however that linear super position doesn't hold as it holds in Maxwell phase. Gauge current is proportional to wave vector and its divergence vanishes as a consequence. Also cylindrically symmetric solutions for which the transverse coordinate is replaced with the radial coordinate  $\rho = \sqrt{m_1^2 + m_2^2}$  are possible. In fact, v can be any function of the coordinates  $m^1, m^2$  transversal to the light like vector p.

Boundary conditions on the boundaries of the massless extremal are satisfied provided the normal component of the energy momentum tensor vanishes. Since energy momentum tensor is of the form  $T^{\alpha\beta} \propto p^{\alpha}p^{\beta}$  the conditions  $T^{n\beta} = 0$  are satisfied if the  $M^4$  projection of the boundary is given by the equations of form

$$H(p \cdot m, \varepsilon \cdot m, \varepsilon_1 \cdot m) = 0 , \qquad \varepsilon_1 \cdot p = 0 , \quad \varepsilon \cdot \varepsilon_1 = 0 .$$
(52)

where H is arbitrary function of its arguments. Recall that for  $M^4$  type extremals the boundary conditions are also satisfied if Kähler field vanishes identically on the boundary.

The following argument suggests that there are not very many manners to satisfy boundary conditions in case of  $M^4$  type extremals. The boundary conditions, when applied to  $M^4$  coordinates imply the vanishing of the normal component of energy momentum tensor. Using coordinates, where energy momentum tensor is diagonal, the requirement boils down to the condition that at least one of the eigen values of  $T^{\alpha\beta}$  vanishes so that the determinant  $det(T^{\alpha\beta})$  must vanish on the boundary: this condition defines 3-dimensional surface in  $X^4$ . In addition, the normal of this surface must have same direction as the eigen vector associated with the vanishing eigen value: this means that three additional conditions must be satisfied and this is in general true in single point only. The boundary conditions in  $CP_2$  coordinates are satisfied provided that the conditions

$$J^{n\beta}J^k_{\ l}\partial_\beta s^l = 0$$

are satisfied. The identical vanishing of the normal components of Kähler electric and magnetic fields on the boundary of massless extremal property provides a manner to satisfy all boundary conditions but it is not clear whether there are any other manners to satisfy them.

The characteristic feature of the massless extremals is that in general the Kähler gauge current is non-vanishing. In ordinary Maxwell electrodynamics this is not possible. This means that these extremals are accompanied by vacuum current, which contains in general case both weak and electromagnetic terms as well as color part.

A possible interpretation of the solution is as the exterior space-time to a topologically condensed particle with vanishing mass described by massless  $CP_2$  type extremal, say photon or neutrino. In general the surfaces in question have boundaries since the coordinates  $s^k$  are are bounded: this is in accordance with the general ideas about topological condensation. The fact that massless plane wave is associated with  $CP_2$  type extremal combines neatly the wave and particle aspects at geometrical level.

The fractal hierarchy of space-time sheets implies that massless extremals should interesting also in long length scales. The presence of a light like electromagnetic vacuum current implies the generation of coherent photons and also coherent gravitons are generated since the Einstein tensor is also non-vanishing and light like (proportional to  $k^{\alpha}k^{\beta}$ ). Massless extremals play an important role in the TGD based model of bio-system as a macroscopic quantum system. The possibility of vacuum currents is what makes possible the generation of the highly desired coherent photon states.

# 4.5 Generalization of the solution ansatz defining massless extremals (MEs)

The solution ansatz for MEs has developed gradually to an increasingly general form and the following formulation is the most general one achieved hitherto. Rather remarkably, it rather closely resembles the solution ansatz for the  $CP_2$  type extremals and has direct interpretation in terms of geometric optics. Equally remarkable is that the latest generalization based on the introduction of the local light cone coordinates was inspired by quantum holography principle.

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#### 4.5.1 Local light cone coordinates

The solution involves a decomposition of  $M_+^4$  tangent space localizing the decomposition of Minkowski space to an orthogonal direct sum  $M^2 \oplus E^2$  defined by light-like wave vector and polarization vector orthogonal to it. This decomposition defines what might be called local light cone coordinates.

1. Denote by  $m^i$  the linear Minkowski coordinates of  $M^4$ . Let  $(S^+, S^-, E^1, E^2)$  denote local coordinates of  $M^4_+$  defining a *local* decomposition of the tangent space  $M^4$  of  $M^4_+$  into a direct *orthogonal* sum  $M^4 = M^2 \oplus E^2$  of spaces  $M^2$  and  $E^2$ . This decomposition has interpretation

in terms of the longitudinal and transversal degrees of freedom defined by local light-like four-velocities  $v_{\pm} = \nabla S_{\pm}$  and polarization vectors  $\epsilon_i = \nabla E^i$  assignable to light ray.

2. With these assumptions the coordinates  $(S_{\pm}, E^i)$  define local light cone coordinates with the metric element having the form

$$ds^{2} = 2g_{+-}dS^{+}dS^{-} + g_{11}(dE^{1})^{2} + g_{22}(dE^{2})^{2} .$$
(53)

If complex coordinates are used in transversal degrees of freedom one has  $g_{11} = g_{22}$ .

3. This family of light cone coordinates is not the most general family since longitudinal and transversal spaces are orthogonal. One can also consider light-cone coordinates for which one non-diagonal component, say  $m_{1+}$ , is non-vanishing if the solution ansatz is such that longitudinal and transversal spaces are orthogonal for the induced metric.

#### 4.5.2 A conformally invariant family of local light cone coordinates

The simplest solutions to the equations defining local light cone coordinates are of form  $S_{\pm} = k \cdot m$ giving as a special case  $S_{\pm} = m^0 \pm m^3$ . For more general solutions of from

$$S_{\pm} = m^0 \pm f(m^1, m^2, m^3)$$
,  $(\nabla_3 f)^2 = 1$ ,

where f is an otherwise arbitrary function, this relationship reads as

$$S^+ + S^- = 2m^0$$

This condition defines a natural rest frame. One can integrate f from its initial data at some two-dimensional f = constant surface and solution describes curvilinear light rays emanating from this surface and orthogonal to it. The flow velocity field  $\overline{v} = \nabla f$  is irrotational so that closed flow lines are not possible in a connected region of space and the condition  $\overline{v}^2 = 1$  excludes also closed flow line configuration with singularity at origin such as  $v = 1/\rho$  rotational flow around axis.

One can identify  $E^2$  as a local tangent space spanned by polarization vectors and orthogonal to the flow lines of the velocity field  $\overline{v} = \nabla f(m^1, m^2, m^3)$ . Since the metric tensor of any 3-dimensional space allows always diagonalization in suitable coordinates, one can always find coordinates  $(E^1, E^2)$  such that  $(f, E^1, E^2)$  form orthogonal coordinates for  $m^0 = constant$  hyperplane. Obviously one can select the coordinates  $E^1$  and  $E^2$  in infinitely many manners.

#### 4.5.3 Closer inspection of the conditions defining local light cone coordinates

Whether the conformal transforms of the local light cone coordinates  $\{S_{\pm} = m^0 \pm f(m^1, m^2, m^3), E^i\}$  define the only possible compositions  $M^2 \oplus E^2$  with the required properties, remains an open question. The best that one might hope is that any function  $S^+$  defining a family of light-like curves defines a local decomposition  $M^4 = M^2 \oplus E^2$  with required properties.

- 1. Suppose that  $S^+$  and  $S^-$  define light-like vector fields which are not orthogonal (proportional to each other). Suppose that the polarization vector fields  $\epsilon_i = \nabla E^i$  tangential to local  $E^2$  satisfy the conditions  $\epsilon_i \cdot \nabla S^+ = 0$ . One can formally integrate the functions  $E^i$  from these condition since the initial values of  $E^i$  are given at  $m^0 = constant$  slice.
- 2. The solution to the condition  $\nabla S_+ \cdot \epsilon_i = 0$  is determined only modulo the replacement

$$\epsilon_i \to \hat{\epsilon}_i = \epsilon_i + k \nabla S_+ \; ,$$

where k is any function. With the choice

$$k = -\frac{\nabla E^i \cdot \nabla S^-}{\nabla S^+ \cdot \nabla S^-}$$

one can satisfy also the condition  $\hat{\epsilon}_i \cdot \nabla S^- = 0$ .

3. The requirement that also  $\hat{\epsilon}_i$  is gradient is satisfied if the integrability condition

$$k = k(S^+)$$

is satisfied: in this case  $\hat{\epsilon}_i$  is obtained by a gauge transformation from  $\epsilon_i$ . The integrability condition can be regarded as an additional, and obviously very strong, condition for  $S^-$  once  $S^+$  and  $E^i$  are known.

4. The problem boils down to that of finding local momentum and polarization directions defined by the functions  $S^+$ ,  $S^-$  and  $E^1$  and  $E^2$  satisfying the orthogonality and integrability conditions

$$(\nabla S^+)^2 = (\nabla S^-)^2 = 0 , \quad \nabla S^+ \cdot \nabla S^- \neq 0 ,$$
  
$$\nabla S^+ \cdot \nabla E^i = 0 , \qquad \qquad \frac{\nabla E^i \cdot \nabla S^-}{\nabla S^+ \cdot \nabla S^-} = k_i(S^+)$$

The number of integrability conditions is 3+3 (all derivatives of  $k_i$  except the one with respect to  $S^+$  vanish): thus it seems that there are not much hopes of finding a solution unless some discrete symmetry relating  $S^+$  and  $S^-$  eliminates the integrability conditions altogether.

A generalization of the spatial reflection  $f \to -f$  working for the separable Hamilton Jacobi function  $S_{\pm} = m^0 \pm f$  ansatz could relate  $S^+$  and  $S^-$  to each other and trivialize the integrability conditions. The symmetry transformation of  $M_+^4$  must perform the permutation  $S^+ \leftrightarrow S^-$ , preserve the light-likeness property, map  $E^2$  to  $E^2$ , and multiply the inner products between  $M^2$ and  $E^2$  vectors by a mere conformal factor. This encourages the conjecture that all solutions are obtained by conformal transformations from the solutions  $S_{\pm} = m^0 \pm f$ .

#### 4.5.4 General solution ansatz for MEs for given choice of local light cone coordinates

Consider now the general solution ansatz assuming that a local wave-vector-polarization decomposition of  $M^4_{\pm}$  tangent space has been found.

- 1. Let  $E(S^+, E^1, E^2)$  be an arbitrary function of its arguments: the gradient  $\nabla E$  defines at each point of  $E^2$  an  $S^+$ -dependent (and thus time dependent) polarization direction orthogonal to the direction of local wave vector defined by  $\nabla S^+$ . Polarization vector depends on  $E^2$  position only.
- 2. Quite a general family of MEs corresponds to the solution family of the field equations having the general form

$$s^k = f^k(S^+, E) \; ,$$

where  $s^k$  denotes  $CP_2$  coordinates and  $f^k$  is an arbitrary function of  $S^+$  and E. The solution represents a wave propagating with light velocity and having definite  $S^+$  dependent polarization in the direction of  $\nabla E$ . By replacing  $S^+$  with  $S^-$  one obtains a dual solution. Field equations are satisfied because energy momentum tensor and Kähler current are light-like so that all tensor contractions involved with the field equations vanish: the orthogonality of  $M^2$ and  $E^2$  is essential for the light-likeness of energy momentum tensor and Kähler current.

- 3. The simplest solutions of the form  $S_{\pm} = m^0 \pm m^3$ ,  $(E^1, E^2) = (m^1, m^2)$  and correspond to a cylindrical MEs representing waves propagating in the direction of the cylinder axis with light velocity and having polarization which depends on point  $(E^1, E^2)$  and  $S^+$  (and thus time). For these solutions four-momentum is light-like: for more general solutions this cannot be the case. Polarization is in general case time dependent so that both linearly and circularly polarized waves are possible. If  $m^3$  varies in a finite range of length L, then 'free' solution represents geometrically a cylinder of length L moving with a light velocity. Of course, ends could be also anchored to the emitting or absorbing space-time surfaces.
- 4. For the general solution the cylinder is replaced by a three-dimensional family of light like curves and in this case the rectilinear motion of the ends of the cylinder is replaced with a curvilinear motion with light velocity unless the ends are anchored to emitting/absorbing space-time surfaces. The non-rotational character of the velocity flow suggests that the freely moving particle like 3-surface defined by ME cannot remain in a infinite spatial volume. The most general ansatz for MEs should be useful in the intermediate and nearby regions of a radiating object whereas in the far away region radiation solution is excepted to decompose to cylindrical ray like MEs for which the function  $f(m^1, m^2, m^2)$  is a linear function of  $m^i$ .
- 5. One can try to generalize the solution ansatz further by allowing the metric of  $M_{+}^{4}$  to have components of type  $g_{i+}$  or  $g_{i-}$  in the light cone coordinates used. The vanishing of  $T^{11}$ ,  $T^{+1}$ , and  $T^{--}$  is achieved if  $g_{i\pm} = 0$  holds true for the induced metric. For  $s^{k} = s^{k}(S^{+}, E^{1})$ ansatz neither  $g_{2\pm}$  nor  $g_{1-}$  is affected by the imbedding so that these components of the metric must vanish for the Hamilton Jacobi structure:

$$ds^{2} = 2g_{+-}dS^{+}dS^{-} + 2g_{1+}dE^{1}dS^{+} + g_{11}(dE^{1})^{2} + g_{22}(dE^{2})^{2} .$$
(54)

 $g_{1+} = 0$  can be achieved by an additional condition

$$m_{1+} = s_{kl}\partial_1 s^k \partial_+ s^k av{55}$$

The diagonalization of the metric seems to be a general aspect of absolute minima. The absence of metric correlations between space-time degrees of freedom for asymptotic selforganization patterns is somewhat analogous to the minimization of non-bound entanglement in the final state of the quantum jump.

#### 4.5.5 Are the boundaries of space-time sheets quite generally light like surfaces with Hamilton Jacobi structure?

Quantum holography principle naturally generalizes to an approximate principle expected to hold true also in non-cosmological length and time scales.

1. The most general ansatz for topological light rays or massless extremals (MEs) inspired by the quantum holographic thinking relies on the introduction of the notion of local light cone coordinates  $S_+, S_-, E_1, E_2$ . The gradients  $\nabla S_+$  and  $\nabla S_-$  define two light like directions just like Hamilton Jacobi functions define the direction of propagation of wave in geometric optics. The two polarization vector fields  $\nabla E_1$  and  $\nabla E_2$  are orthogonal to the direction of propagation defined by either  $S_+$  or  $S_-$ . Since also  $E_1$  and  $E_2$  can be chosen to be orthogonal, the metric of  $M_+^4$  can be written locally as  $ds^2 = g_{+-}dS_+dS_-+g_{11}dE_1^2+g_{22}dE_2^2$ . In the earlier ansatz  $S_+$  and  $S_-$  where restricted to the variables  $k \cdot m$  and  $\tilde{k} \cdot m$ , where k and  $\tilde{k}$  correspond to light like momentum and its mirror image and m denotes linear  $M^4$  coordinates: these MEs describe cylindrical structures with constant direction of wave propagation expected to be most important in regions faraway from the source of radiation.

- 2. Boundary conditions are satisfied if the 3-dimensional boundaries of MEs have one light like direction ( $S_+$  or  $S_-$  is constant). This means that the boundary of ME has metric dimension d = 2 and is characterized by an infinite-dimensional super-canonical and super-conformal symmetries just like the boundary of the imbedding space  $M_+^4 \times CP_2$ : The boundaries are like moments for mini big bangs (in TGD based fractal cosmology big bang is replaced with a silent whisper amplified to not necessarily so big bang).
- 3. These observations inspire the conjecture that boundary conditions for  $M^4$  like space-time sheets fixed by the variational principle selecting preferred extremals of Kähler action quite generally require that space-time boundaries correspond to light like 3-surfaces with metric dimension equal to d = 2. This does not yet imply that light like surfaces of imbedding space would take the role of the light cone boundary: these light like surface could be seen only as a special case of causal determinants analogous to event horizons.

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